

DARK MATTER THEORY II

REBECCA LEANE

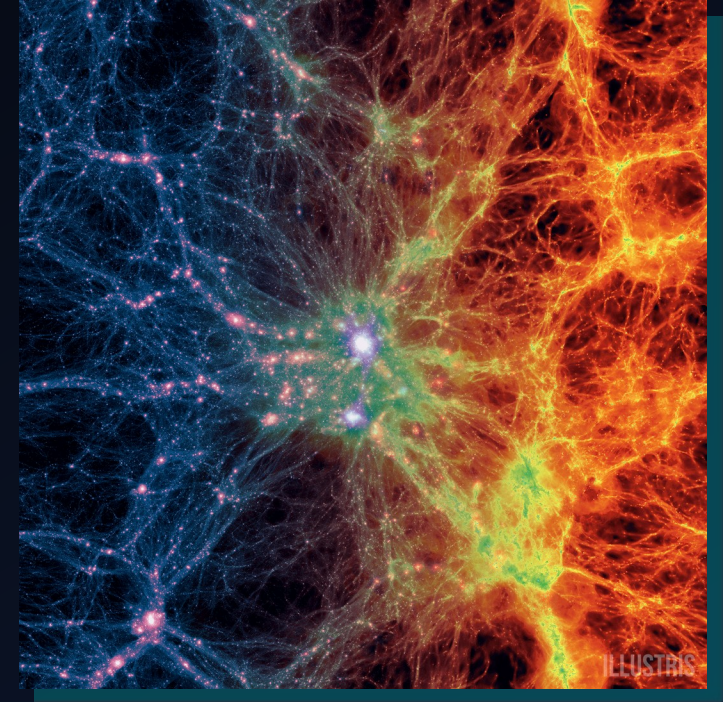
SLAC NATIONAL ACCELERATOR LABORATORY

SLAC SSI LECTURES
AUG 15-16th 2022



Today's Outline

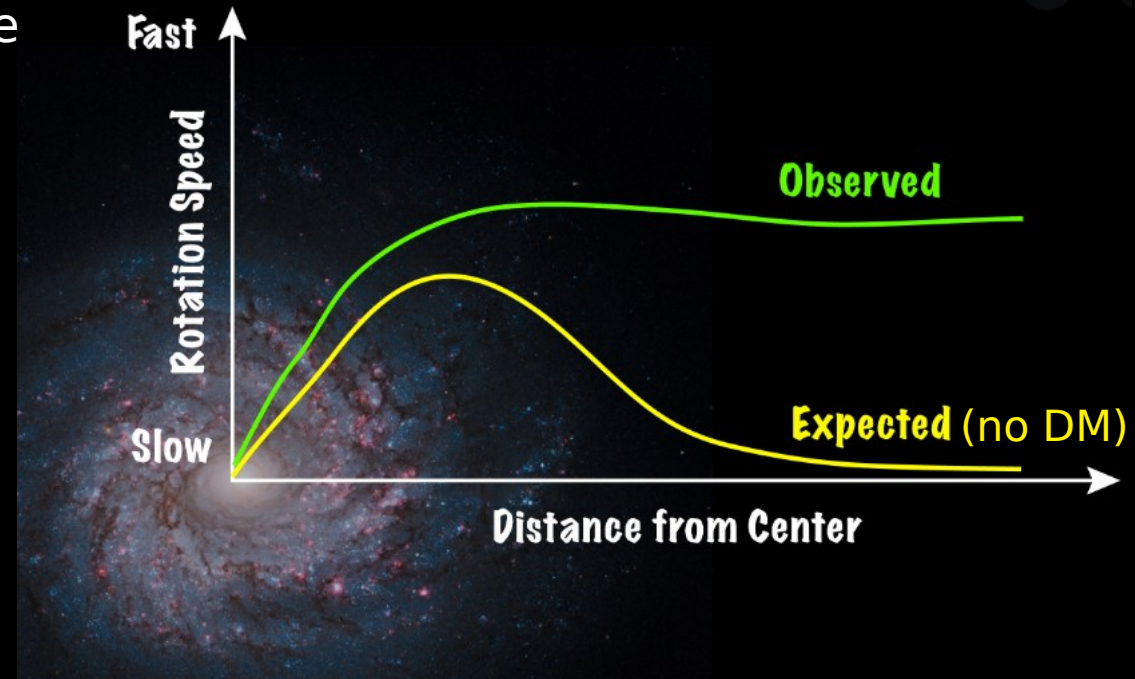
- WIMP recap and status
 - Searches and Constraints
- Newer Production Mechanisms
 - SIMPs, co-SIMPs
 - Freeze-in
- Axions
- MACHOs
 - Primordial Black Holes
- Outlook



Illustris Simulation

Recap: what we know about Dark Matter

- Makes up about 85% of matter in the Universe (~5x more mass than SM)
- Doesn't interact too much with light (or itself)
- Provides large scale structure of the Universe
- Structure tells us non-relativistic
- Forms halos around Galaxies
- Stable on cosmological timescales

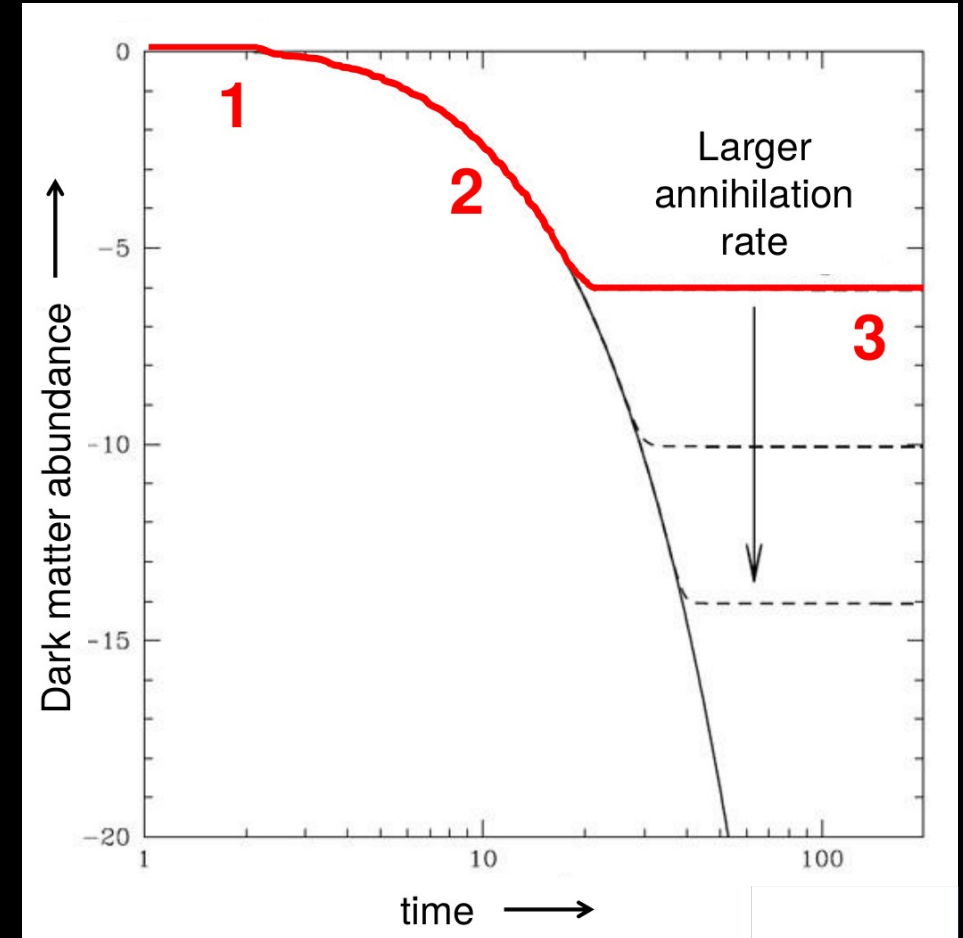


Cannot be explained by any known particles!

Recap: Dark Matter Abundance: WIMPs

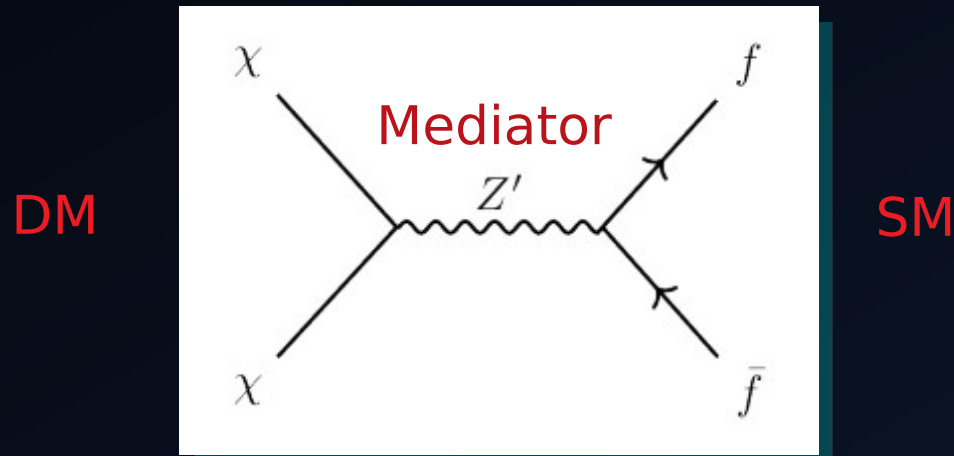
- 1)** Thermal equilibrium:
 $\text{DM} + \text{DM} \Rightarrow \text{visible particles}$
 $\text{Visible particles} \Rightarrow \text{DM} + \text{DM}$
- 2)** Universe cools, only
 $\text{DM} + \text{DM} \Rightarrow \text{visible particles}$
- 3)** Universe expands too fast.
No more annihilations.
DM abundance is set.

Predicts a particular annihilation rate for dark matter.



Recap: Dark Sectors

- Standard Model (SM) mediated WIMP processes, e.g. with Z bosons, extremely constrained
- Need new fields beyond the SM!
- These can mediate interactions between the dark sector, and the SM



Dark Sectors

- Need new fields beyond the SM, consider extensions:
 - New scalar, pseudoscalar, fermion, or vector
- Symmetries of SM restrict how these fields interact with the SM

$$-\lambda_{HS}(S^\dagger S)(H^\dagger H)$$

Higgs Portal

$$\frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

Axion Portal

$$-\frac{\epsilon}{2} B_{\mu\nu} Z'^{\mu\nu}$$

Hypercharge Portal

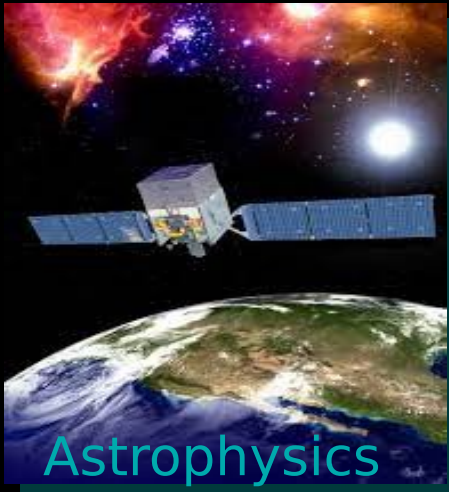
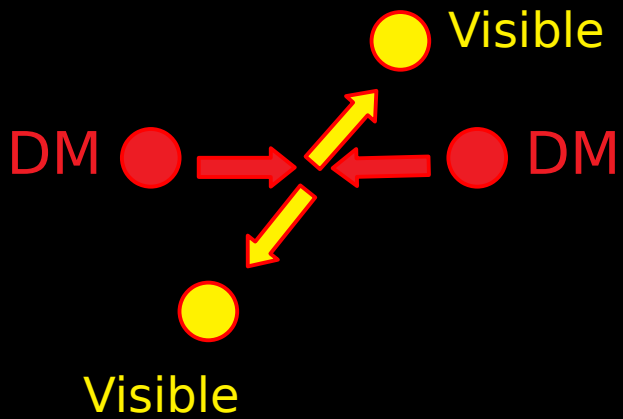
$$y_n L H N$$

Fermion Portal

Generally easiest to parameterize searches by generic masses and couplings

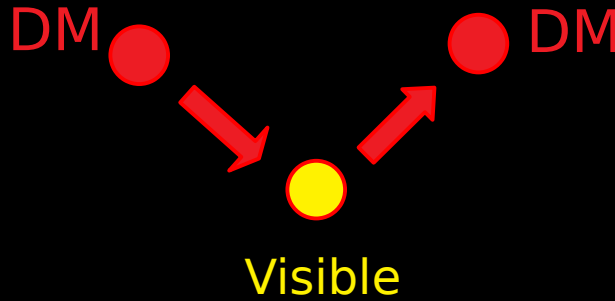
WIMP Dark Matter Search Program

Annihilation



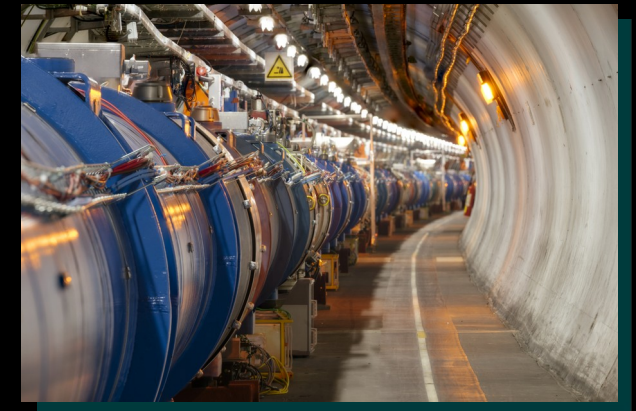
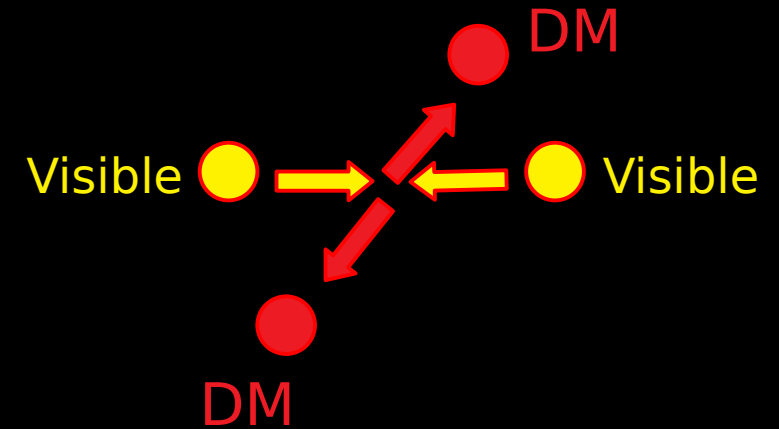
Astrophysics

Scattering



Direct detection
+ Astrophysics

Production

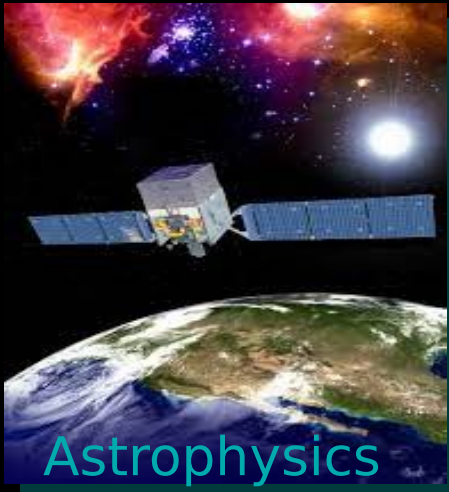
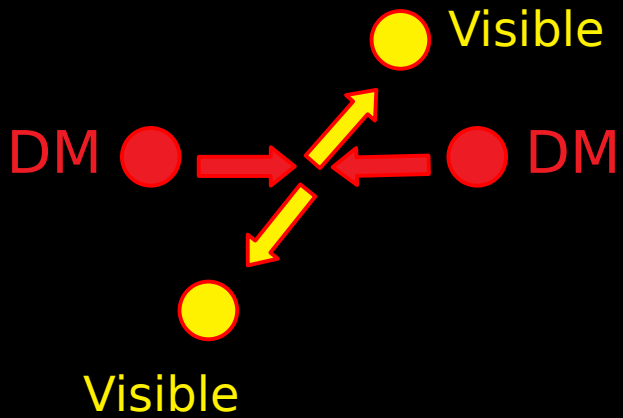


Colliders
+ Astrophysics

WIMP Dark Matter Search Program

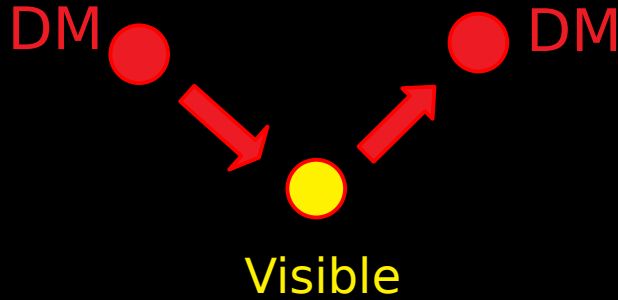
M.E. Mozani's lecture yesterday

Annihilation



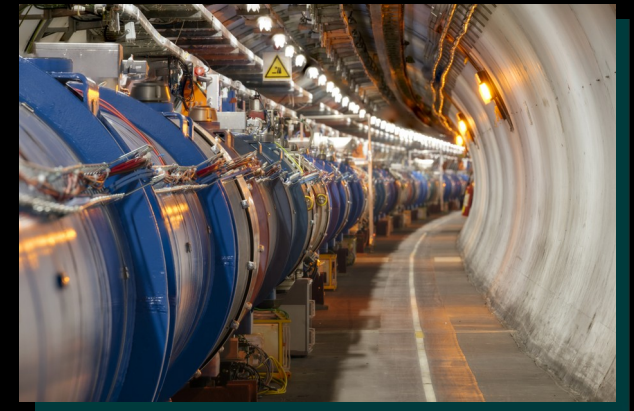
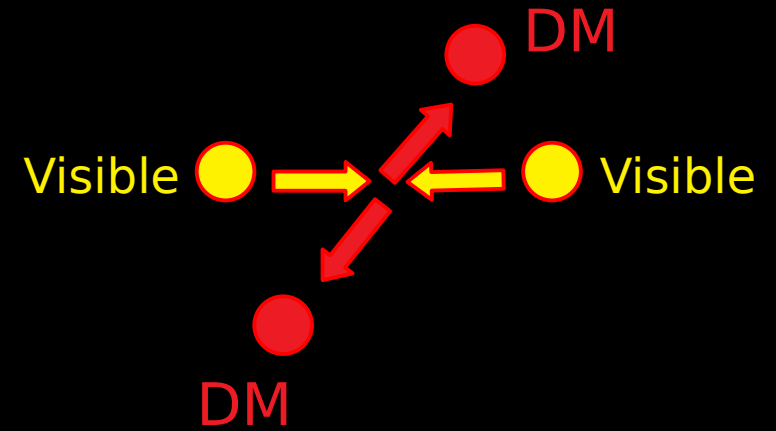
Astrophysics

Scattering



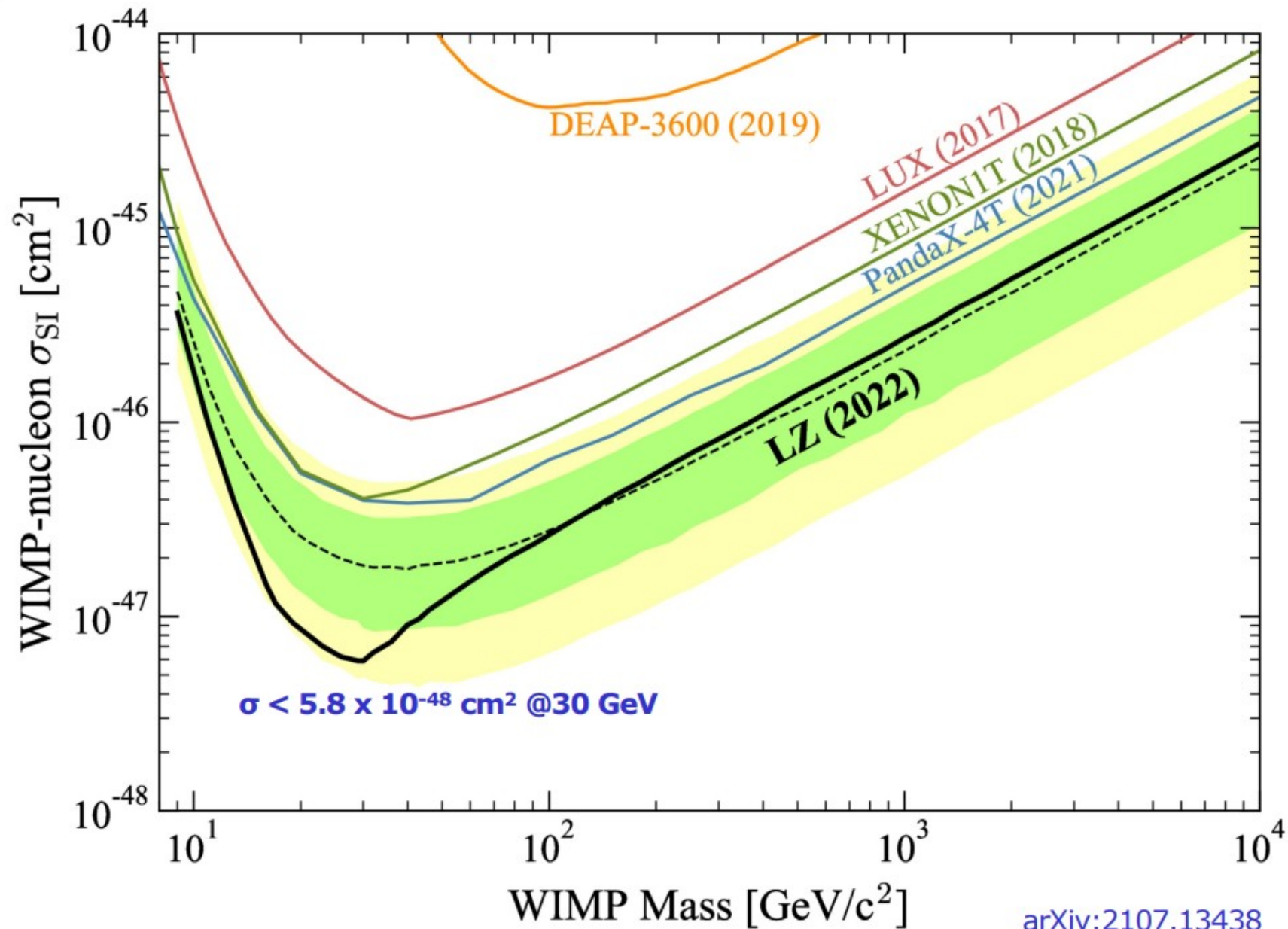
Direct detection
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Production



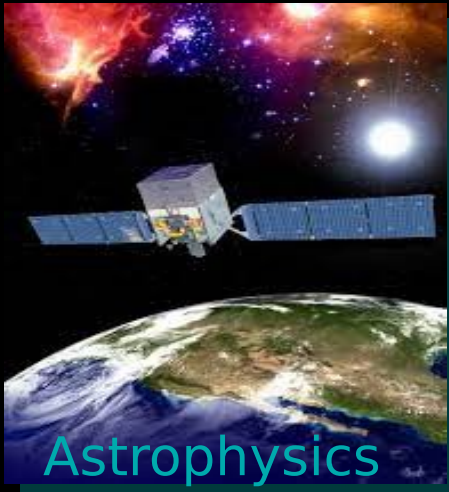
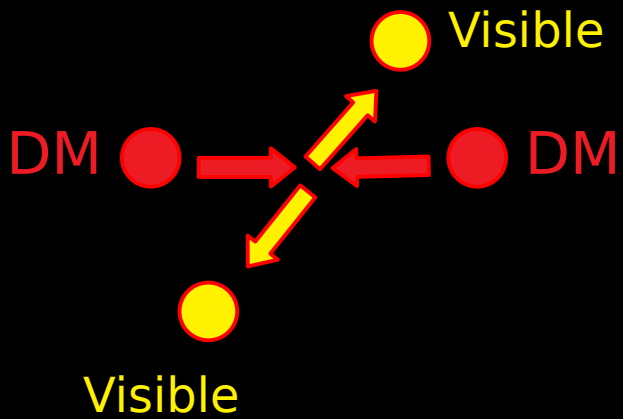
Colliders
+ Astrophysics

World-leading constraint on WIMPs



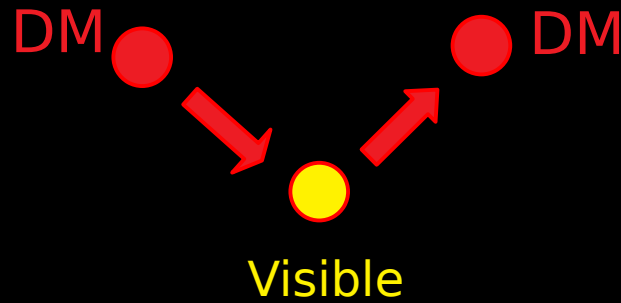
WIMP Dark Matter Search Program

Annihilation



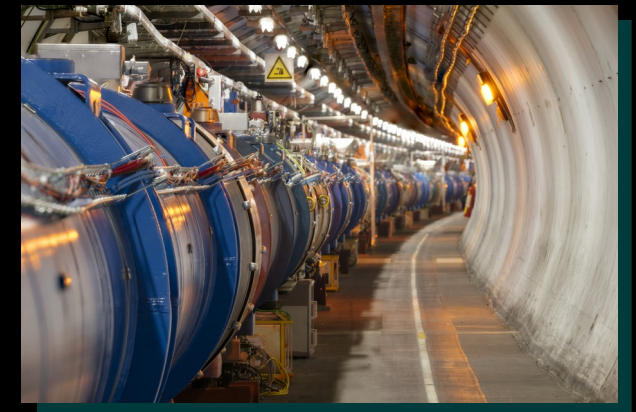
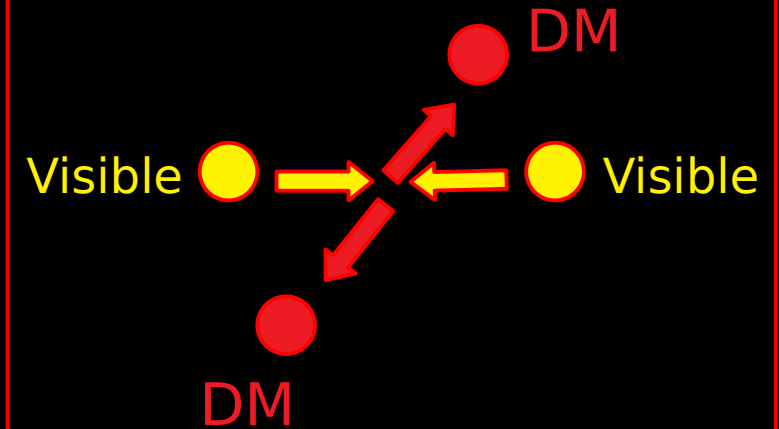
Astrophysics

Scattering



Direct detection
+ Astrophysics

Production



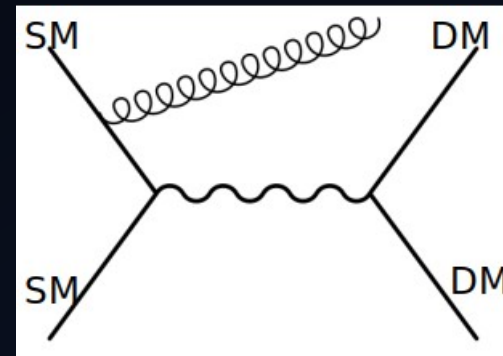
Colliders
+ Astrophysics

Dark matter at colliders



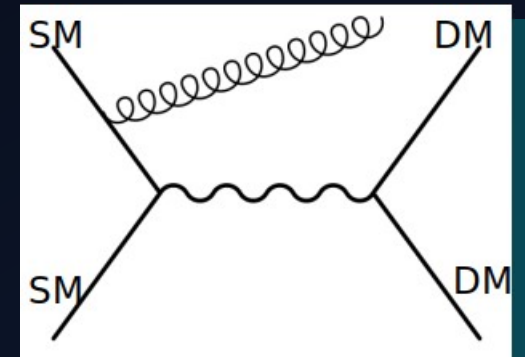
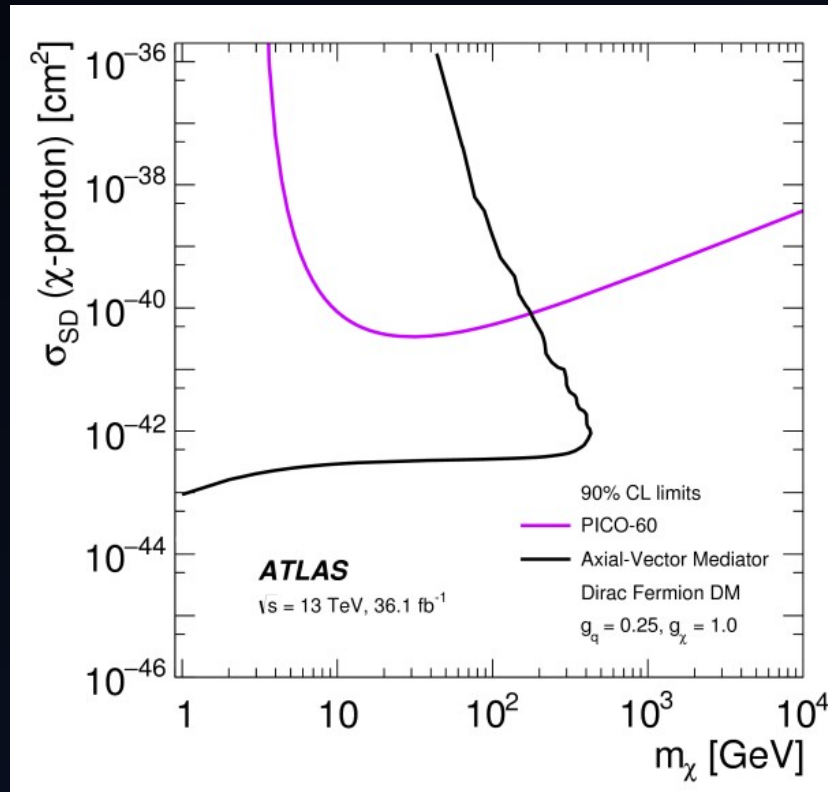
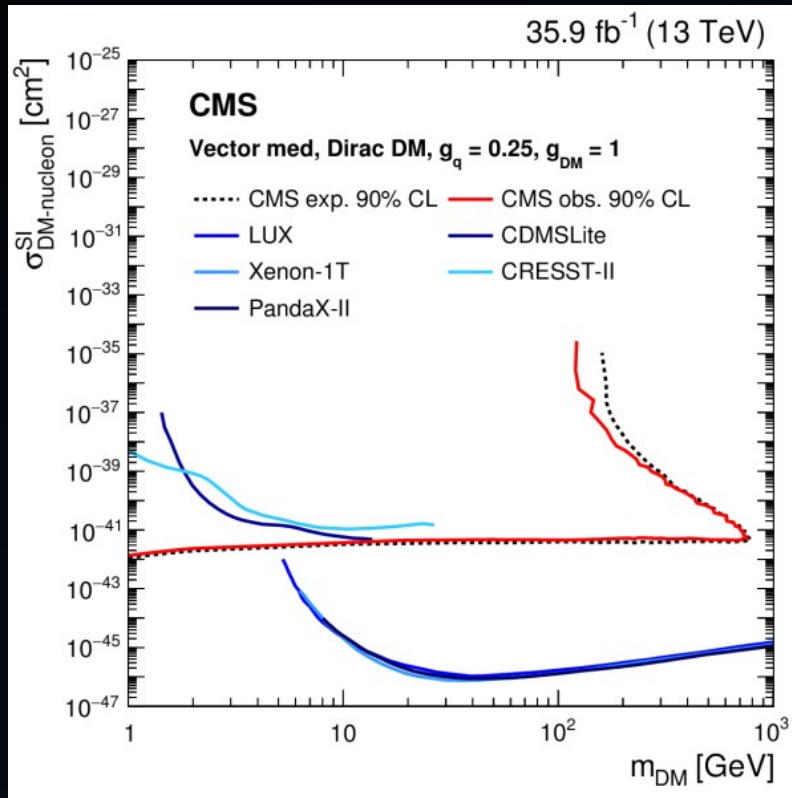
ATLAS

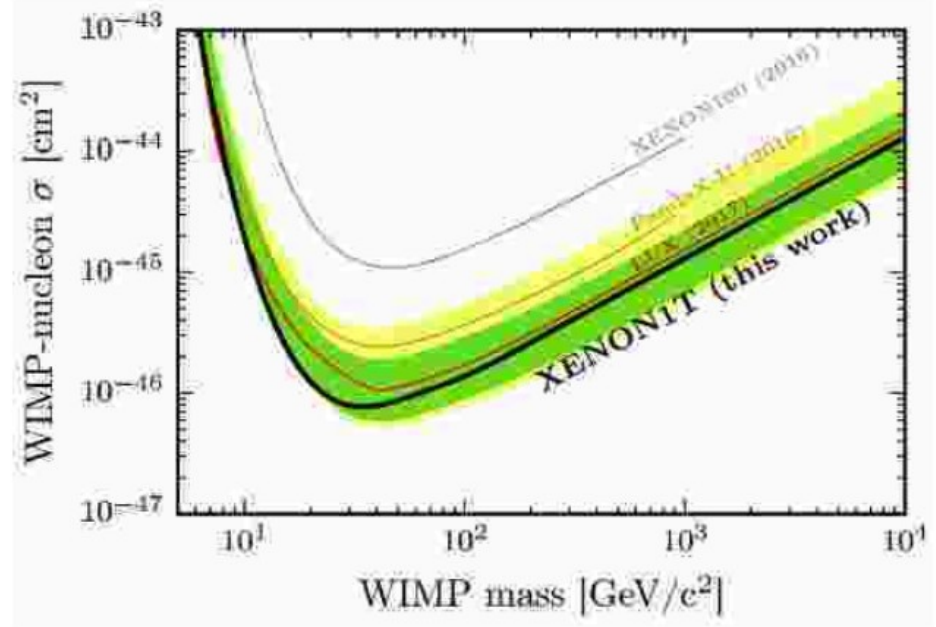
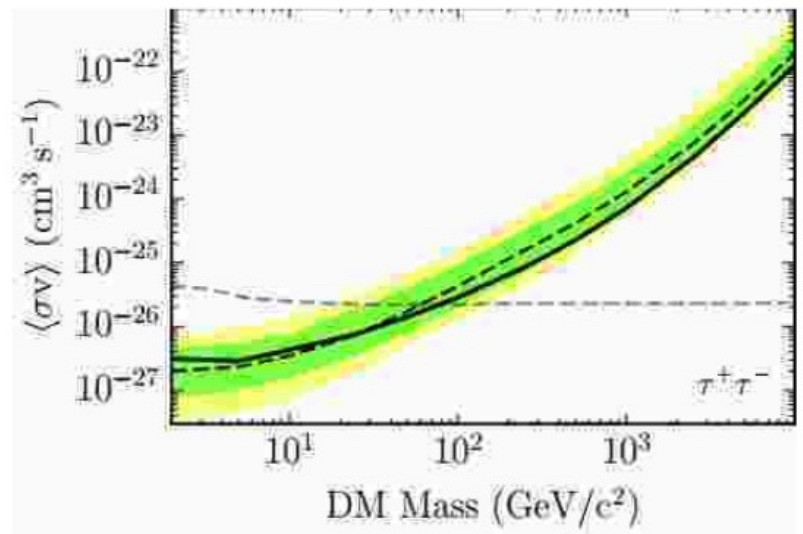
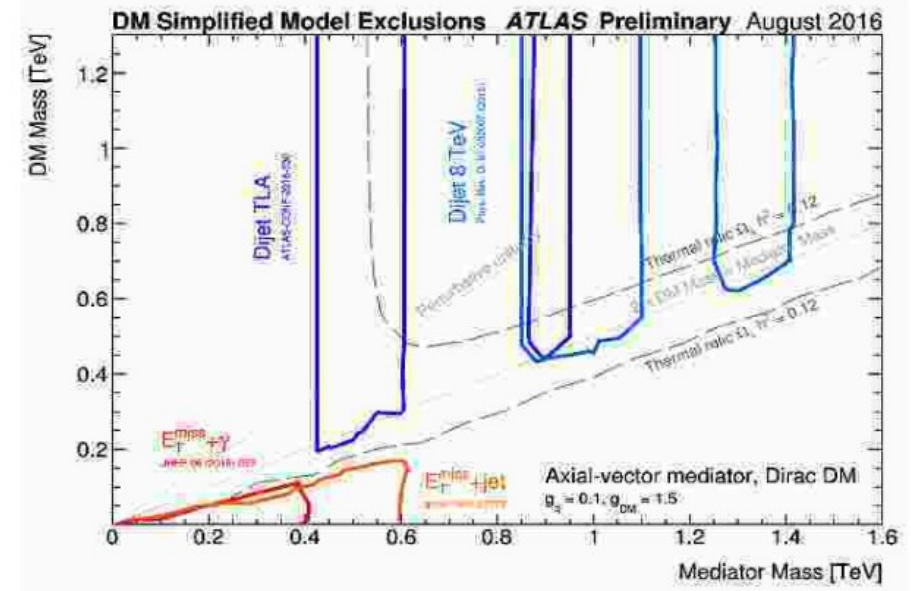
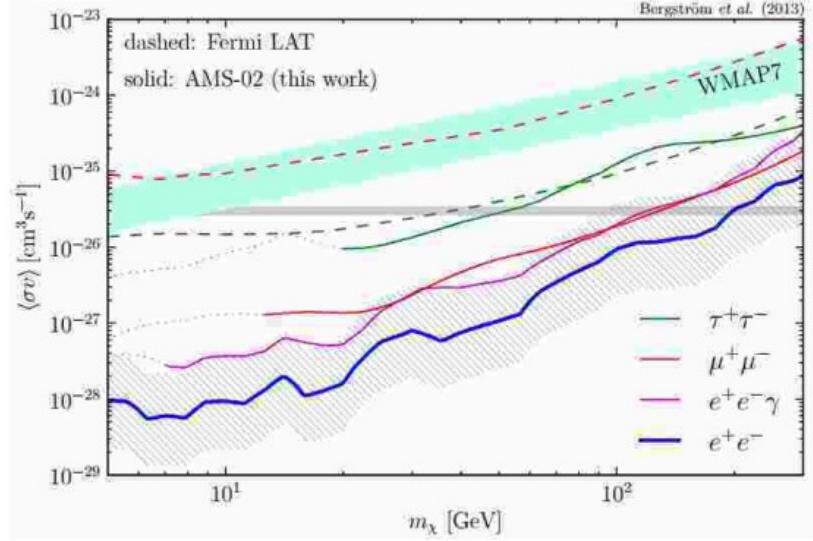
- Missing momentum searches test weak-scale dark matter

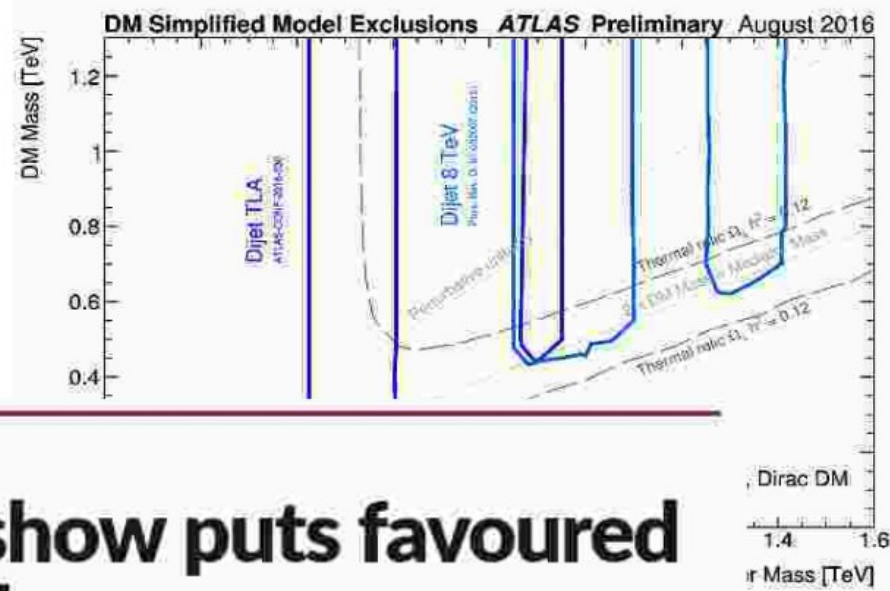
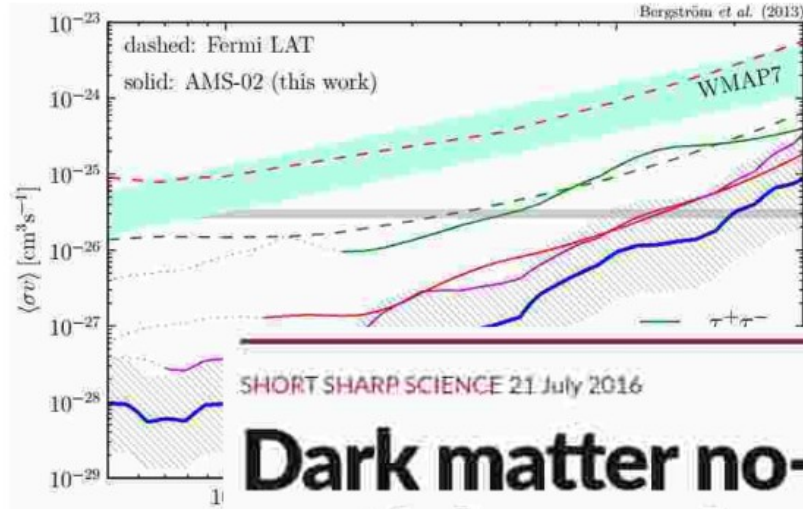


- Strong limits for mediator searches directly as well

Mono-jet searches

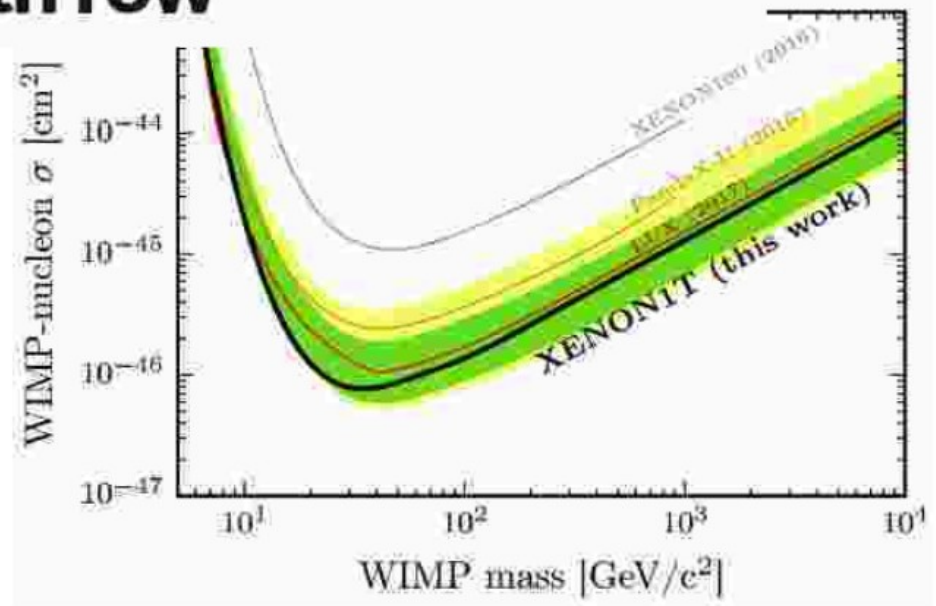
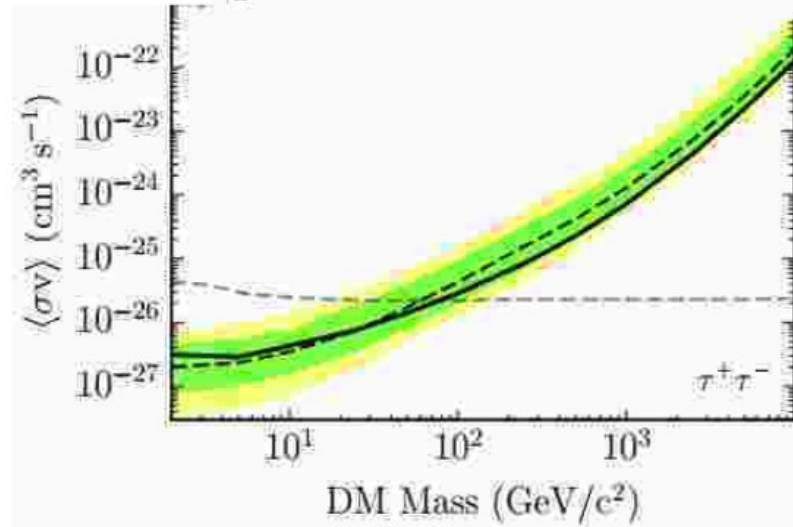


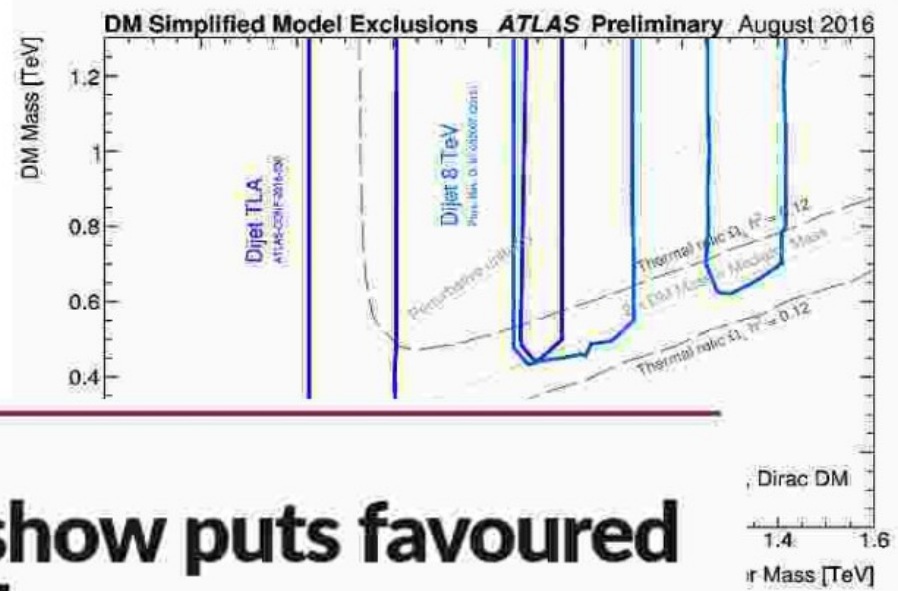
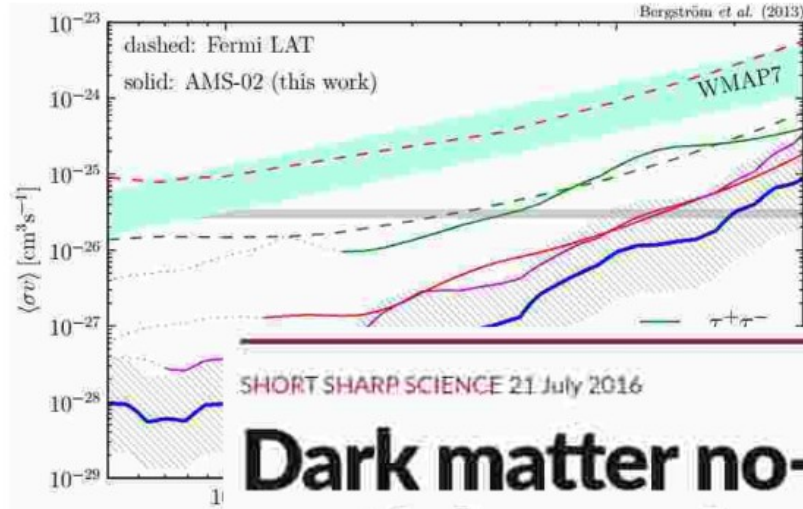




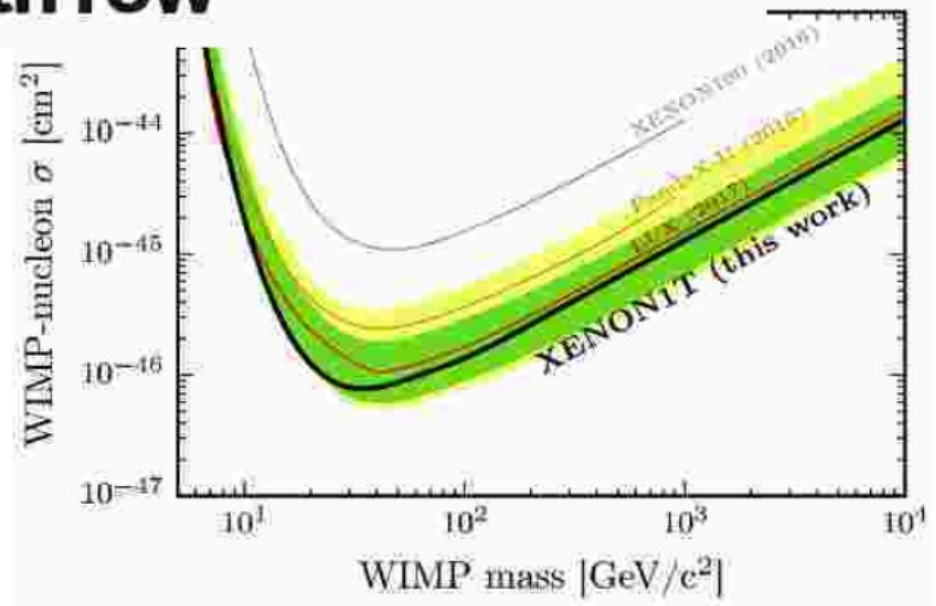
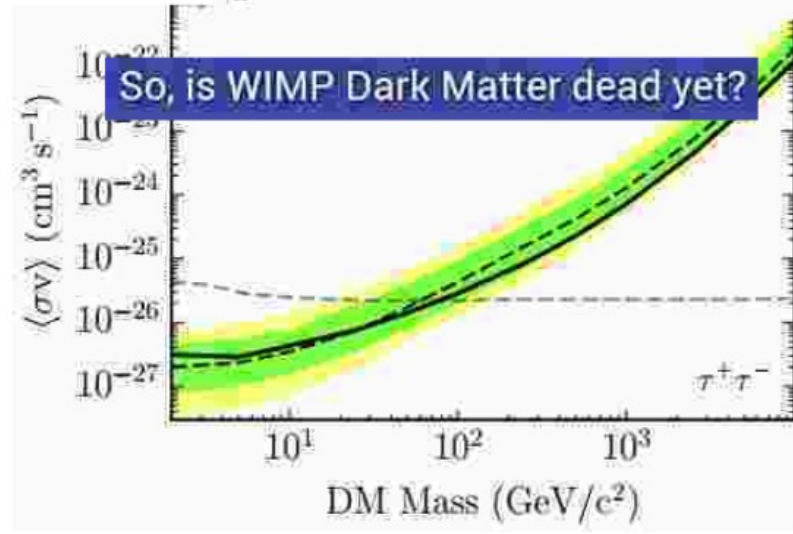
Dark matter no-show puts favoured particles on death row

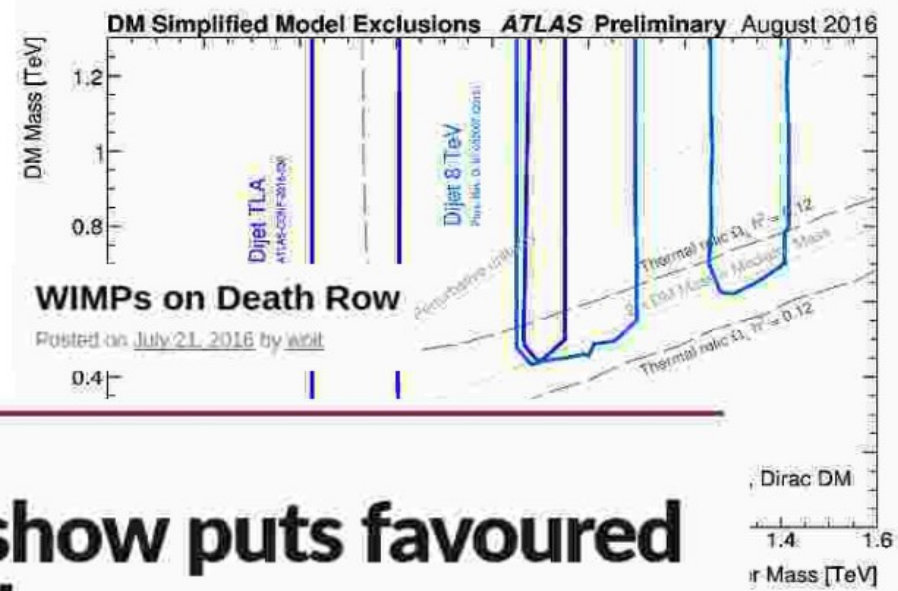
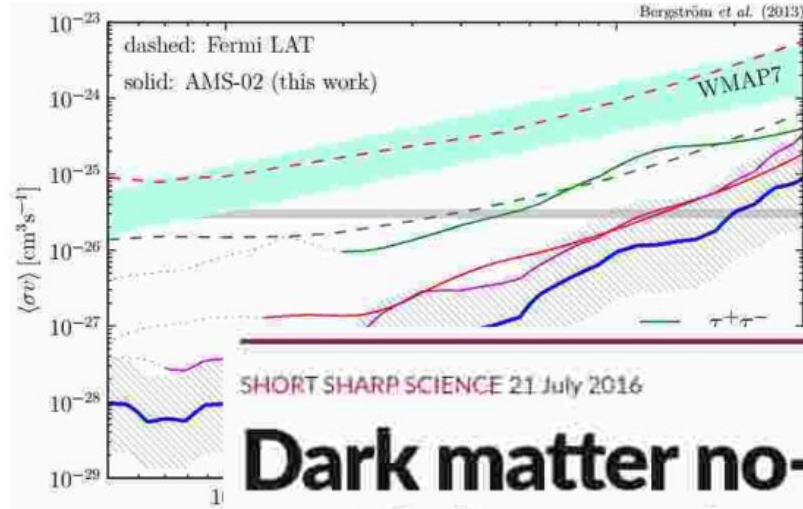
SHORT SHARP SCIENCE 21 July 2016



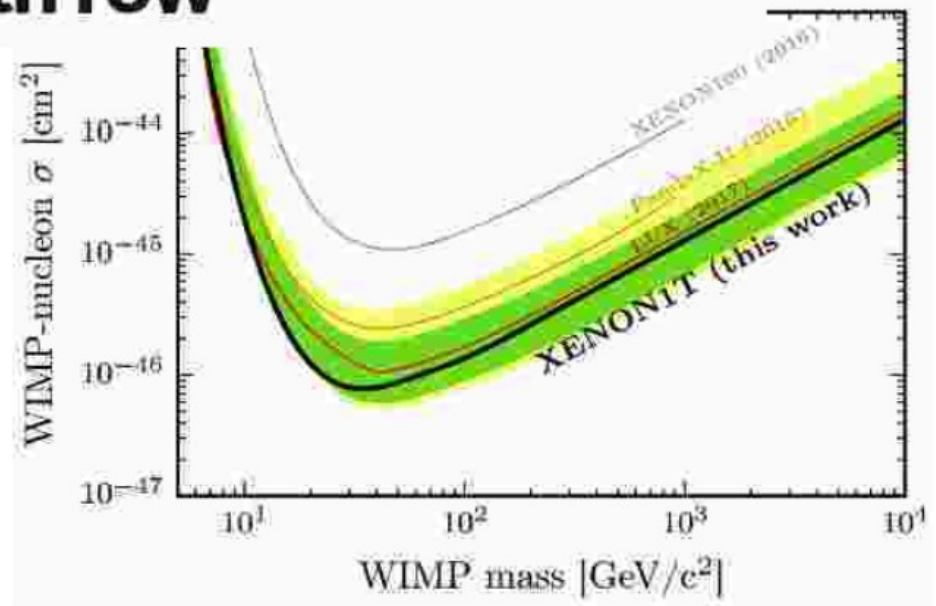
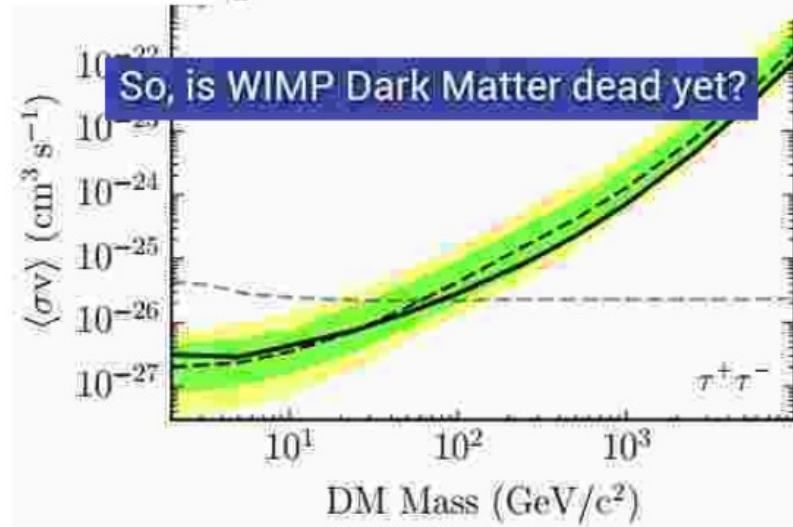


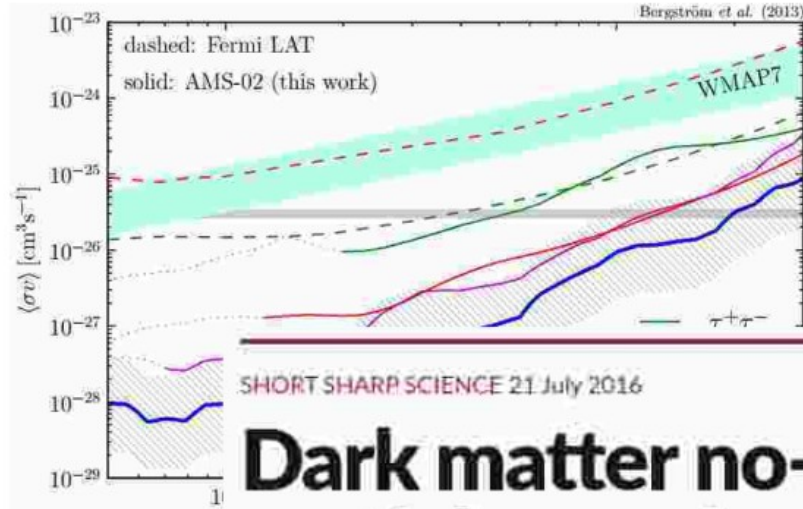
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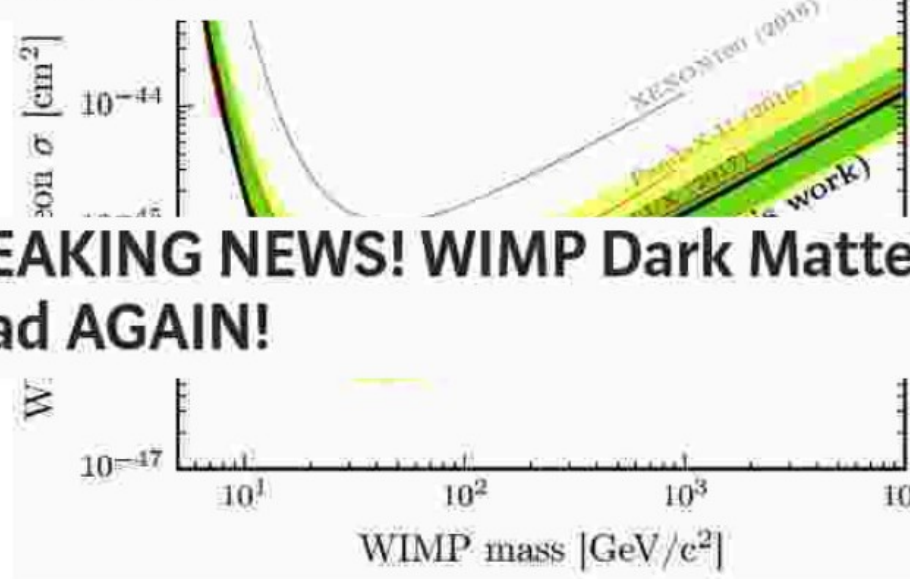
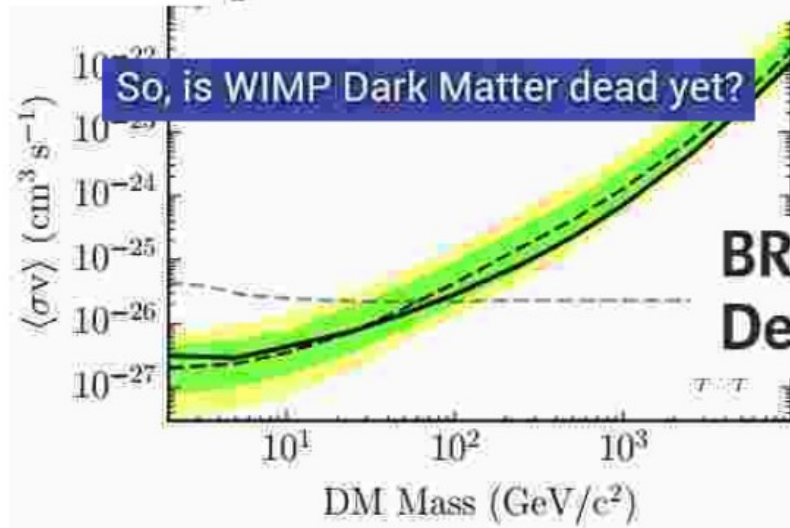
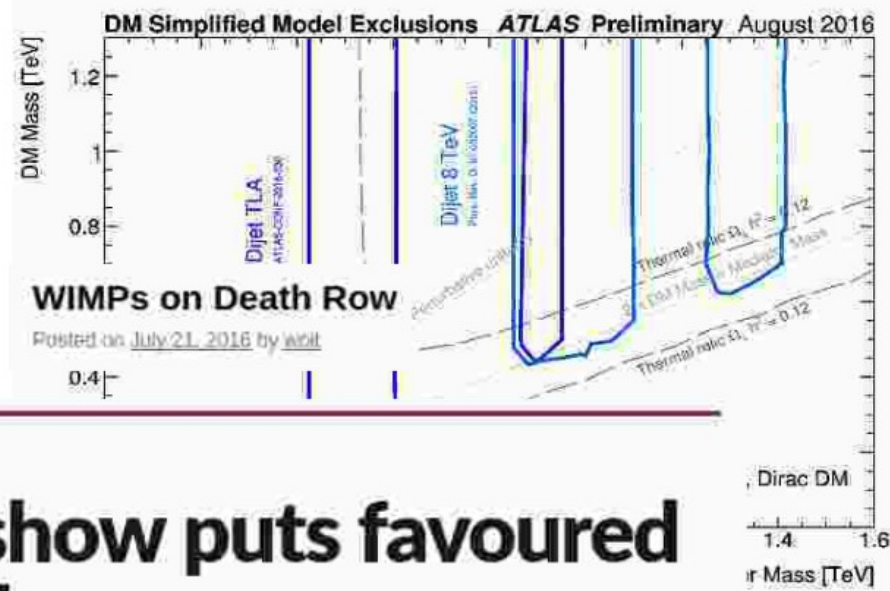
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SHORT SHARP SCIENCE 21 July 2016

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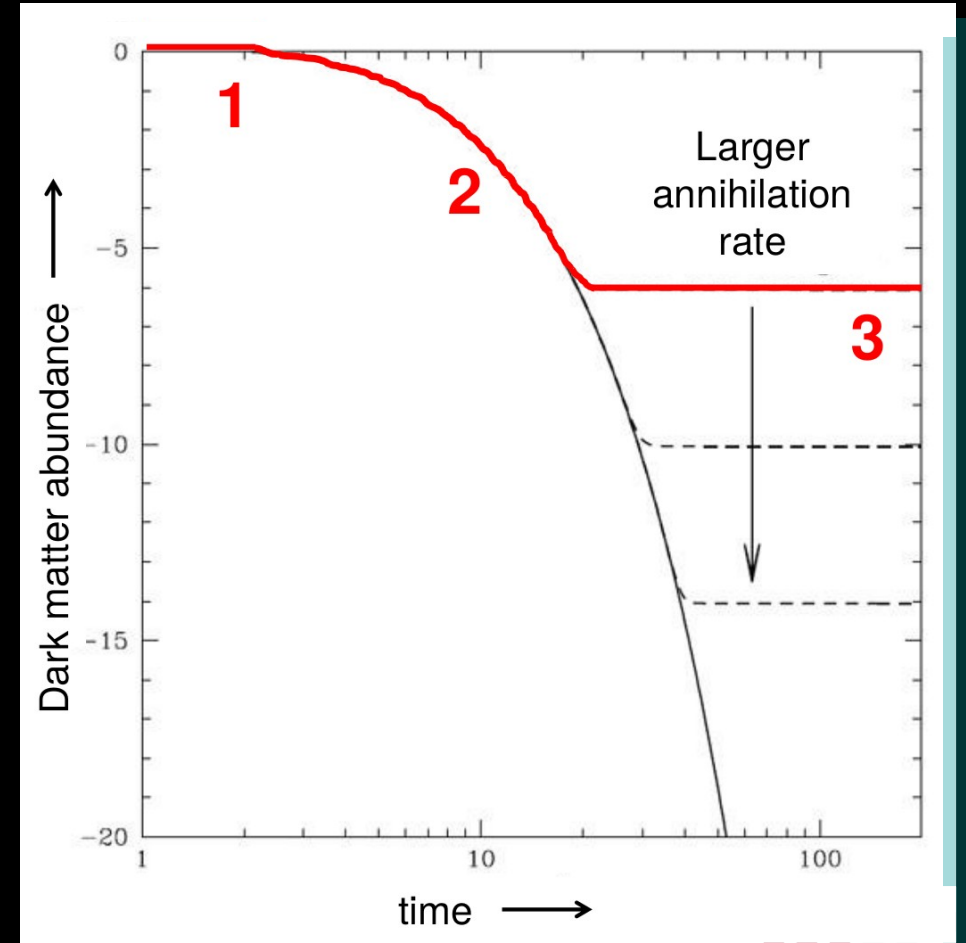
BREAKING NEWS! WIMP Dark Matter Dead AGAIN!

Only annihilation provides the target rate

Scattering (direct detection) and production (colliders):

No well-defined scale because only some aspects of the interaction are being considered.

Those constraints tell you which WIMP models are excluded, but nothing about the remaining possibilities.



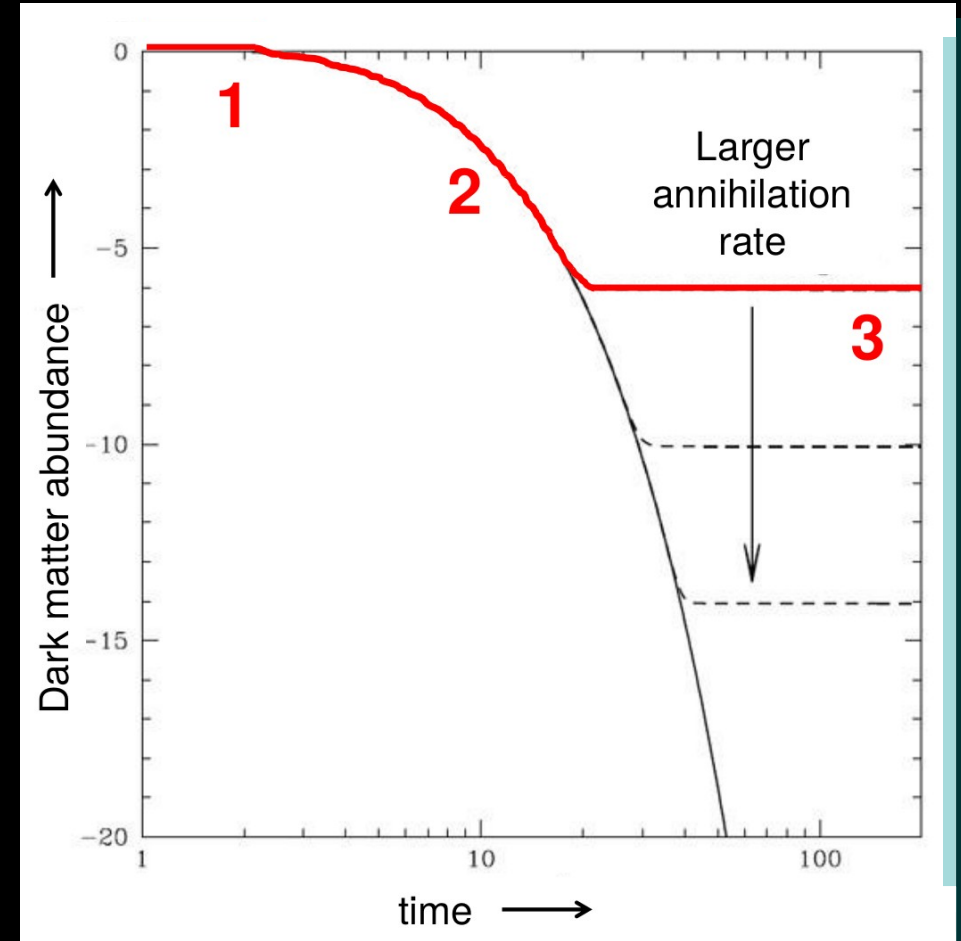
Only annihilation provides the target rate

Scattering (direct detection) and production (colliders):

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Those constraints tell you which WIMP models are excluded, but nothing about the remaining possibilities.

The most decisive way to test thermal WIMPs is through their annihilation products, as this exactly goes to their most fundamental feature: being annihilation relics, which sets a well-defined scale for the total cross section.



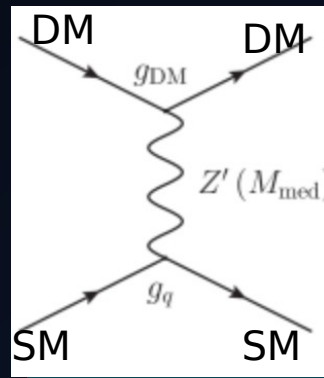
Hidden dark sectors:

- New fields that have **small couplings** to the SM
- Naturally occurs in many model scenarios; symmetries, or Higgs may have mass coupling dependence, etc
- If zero coupling to the SM, can't detect the DM!
- Need to **interact at least very weakly** with the SM to probe at particle experiments

Searching for hidden dark sectors

- Small coupling to SM minimizes constraints from scattering or production
 - Difficult but not impossible to probe!

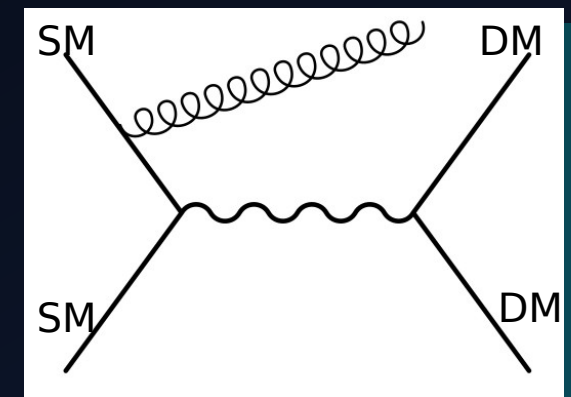
Scattering



Suppressed :(

Rebecca Leane

Production

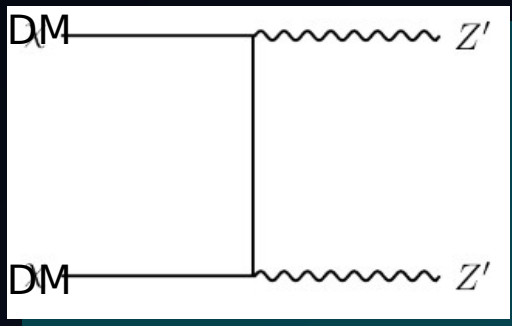


Suppressed :(

Searching for hidden dark sectors

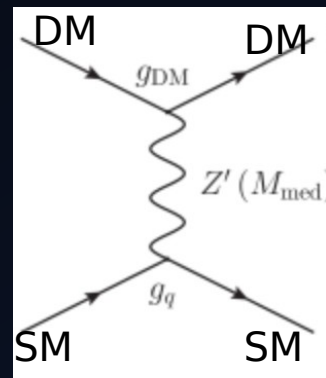
- Small coupling to SM minimizes constraints from scattering or production
 - Difficult but not impossible to probe!
- Large signals still possible for annihilation!
 - Cross section independent of SM coupling if particles produced on shell
 - Only then multiplied by branching fraction

Annihilation



Not suppressed :)

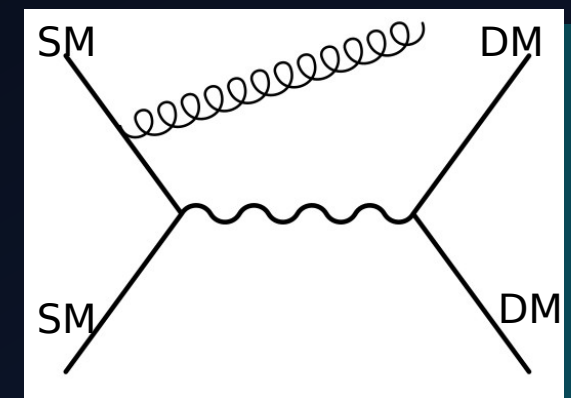
Scattering



Suppressed :(

Rebecca Leane

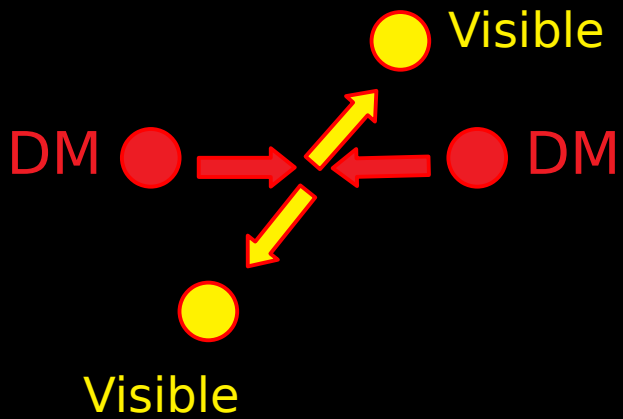
Production



Suppressed :(

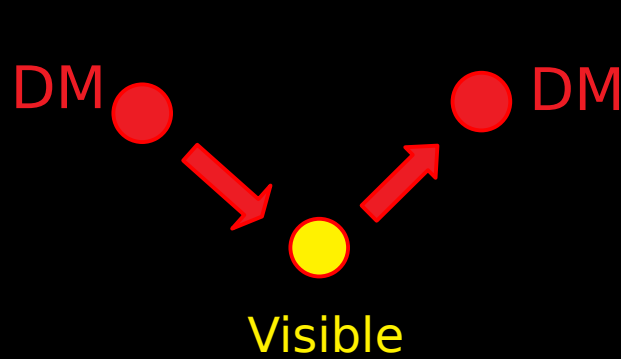
WIMP Dark Matter Search Program

Annihilation



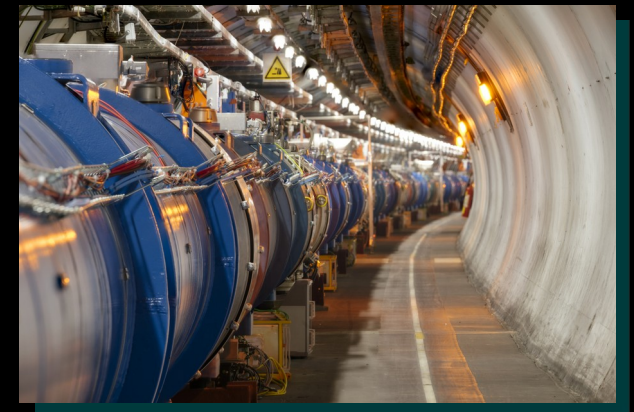
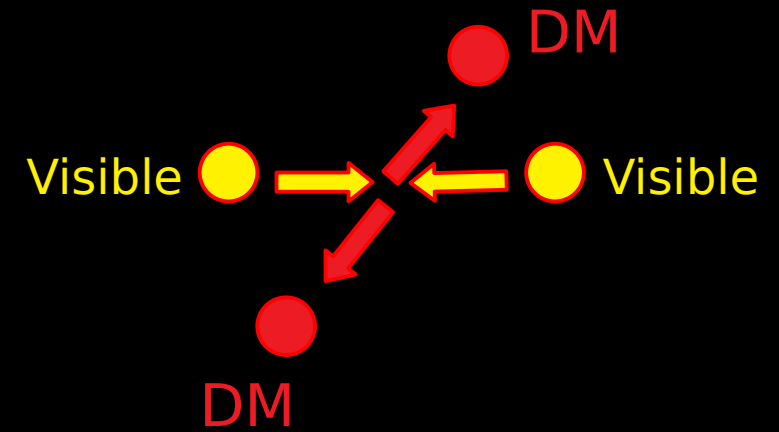
Astrophysics

Scattering



Direct detection
+ Astrophysics

Production



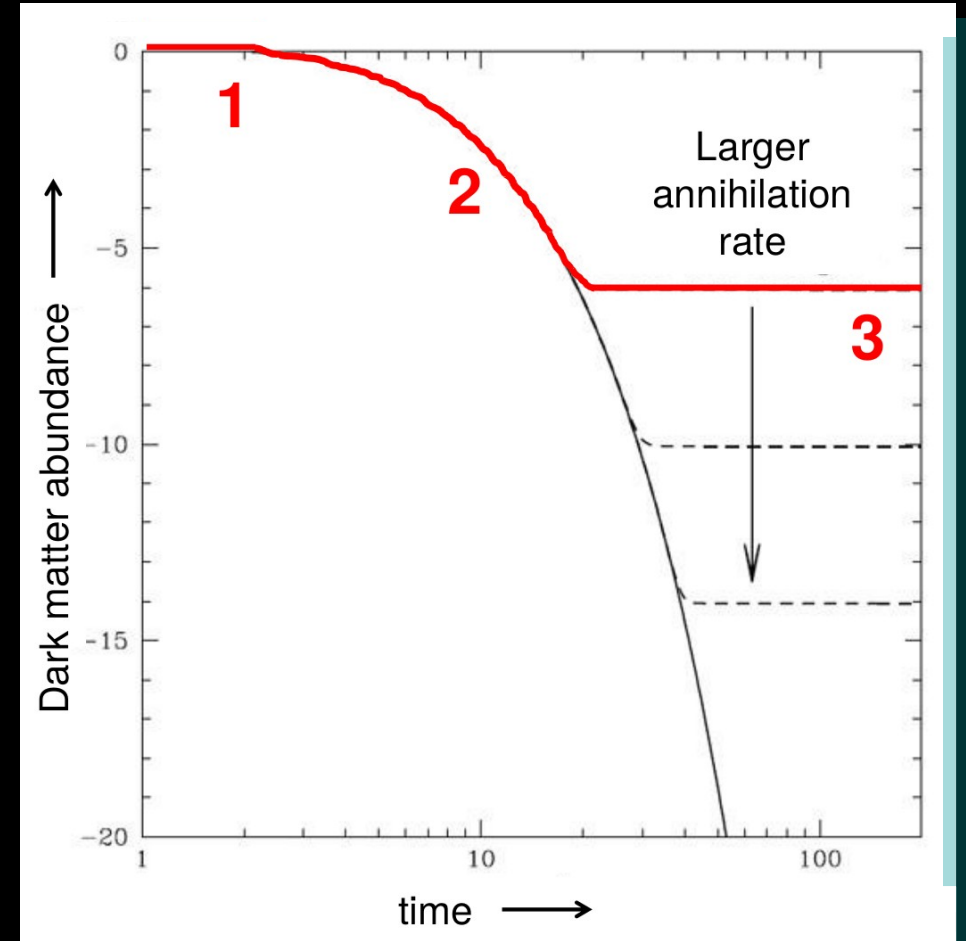
Colliders
+ Astrophysics

INDIRECT DETECTION

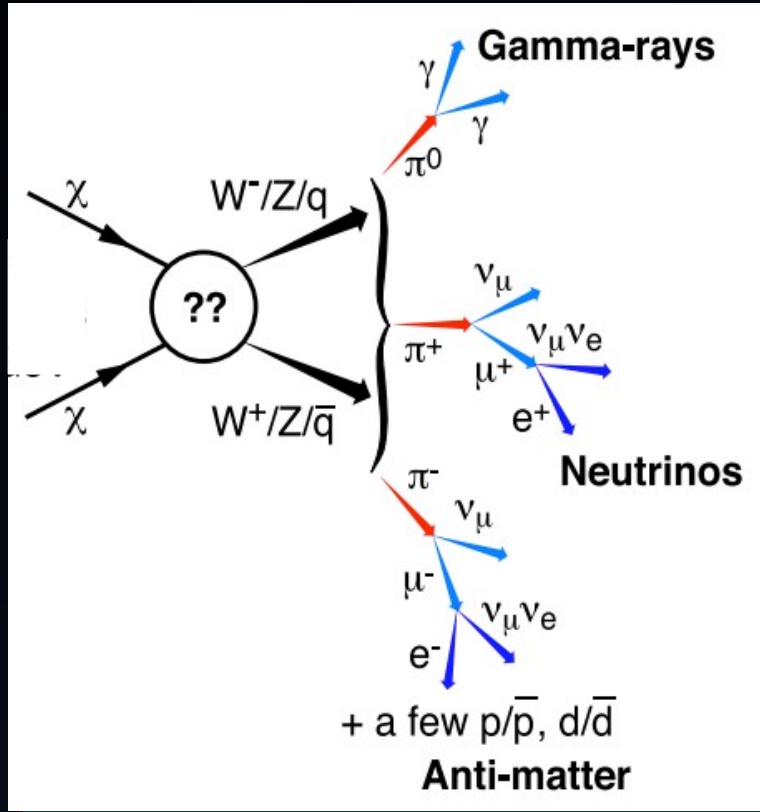
Ingredient #1: DM Interaction Rate

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 $\text{Visible particles} \Rightarrow \text{DM} + \text{DM}$
- 2)** Universe cools, only
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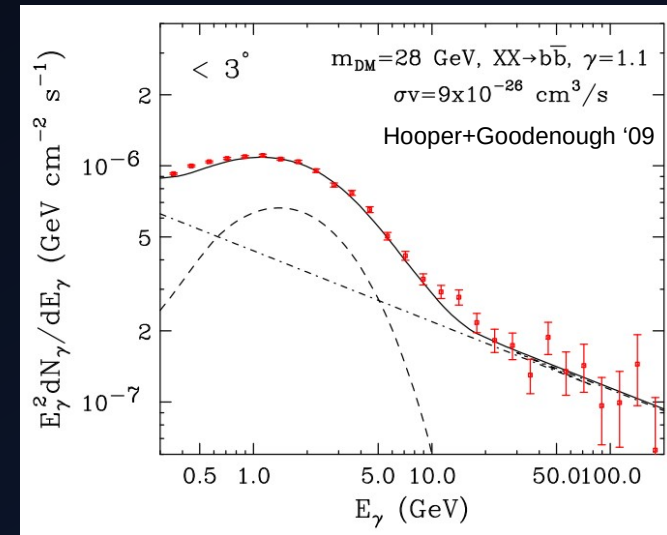


Ingredient #2: Energy Spectrum



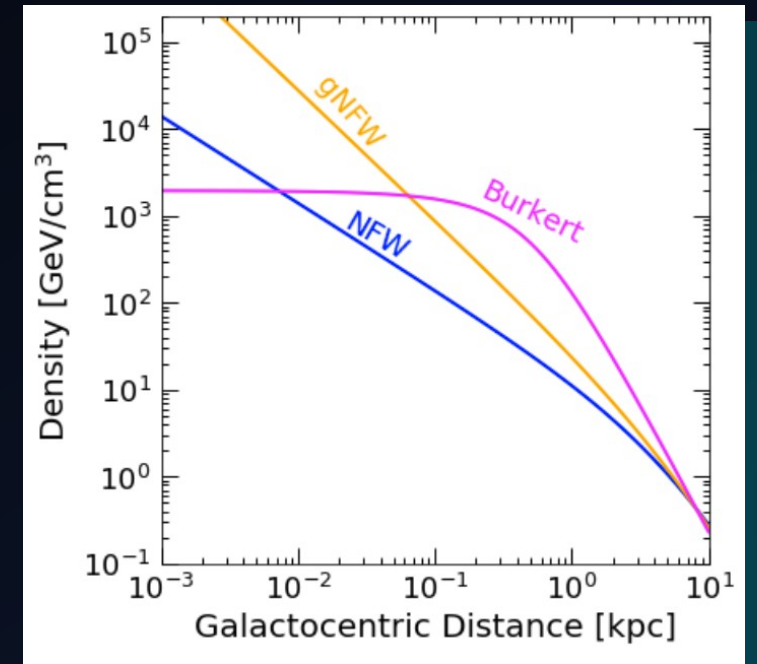
Baltz et al 0806.2911

- Also driven by particle physics model
- Shape depends on:
 - branching ratios to final SM states
 - boosts of particles



Ingredient #3: DM Density+Distribution

- Line of sight integral over DM density
 - J-factor (annihilation)
 - D-factor (decay)
- DM density profiles not well-known
 - large uncertainties



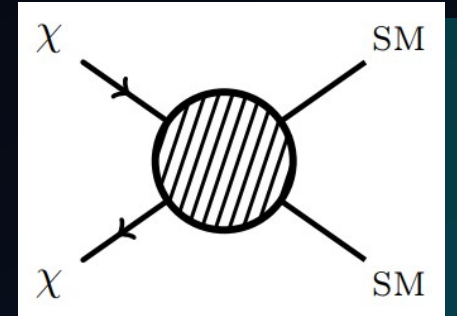
Indirect Detection Ingredients

Particle Physics

Astrophysics

(Neutral particles)

$$\Phi(E, \phi) = \frac{\Gamma}{4\pi m_\chi^a} \frac{dN}{dE} \int \rho[r, (\ell, \phi)]^a d\ell.$$



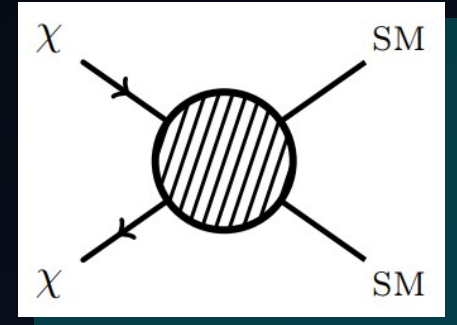
Indirect Detection Ingredients

Particle Physics

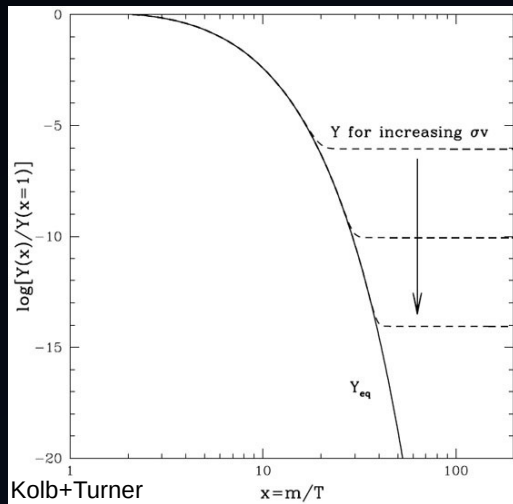
Astrophysics

(Neutral particles)

$$\Phi(E, \phi) \rightarrow \frac{\Gamma}{4\pi m_\chi^a} \frac{dN}{dE} \int \rho[r, (l, \phi)]^a dl.$$



Annihilation cross section



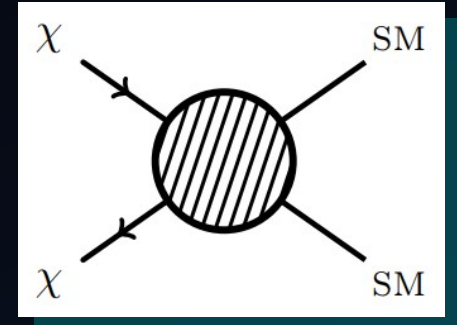
Indirect Detection Ingredients

Particle Physics

Astrophysics

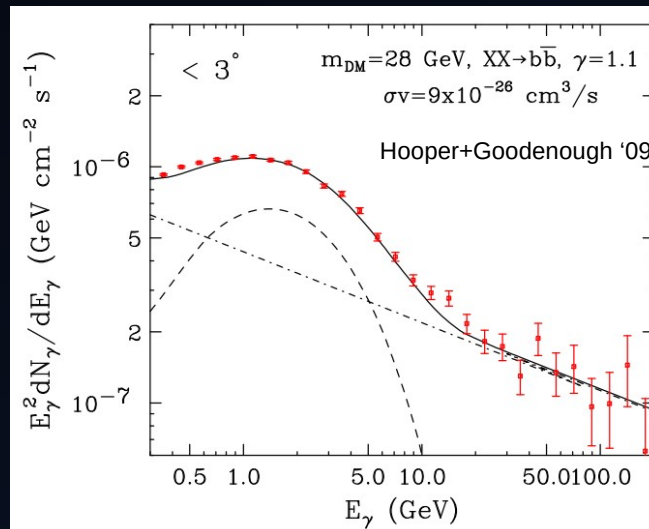
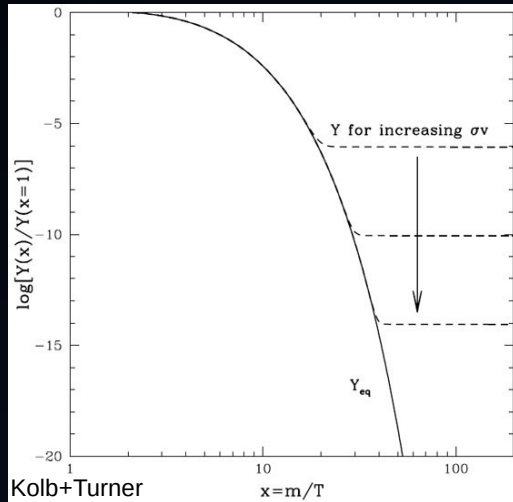
(Neutral particles)

$$\Phi(E, \phi) \propto \frac{\Gamma}{4\pi m_\chi^a} \frac{dN}{dE} \int \rho[r, (\ell, \phi)]^a d\ell.$$



Annihilation cross section

Energy spectrum



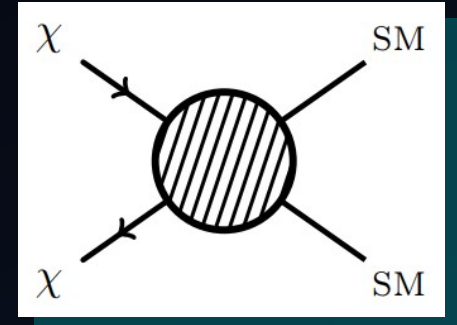
Indirect Detection Ingredients

Particle Physics

Astrophysics

(Neutral particles)

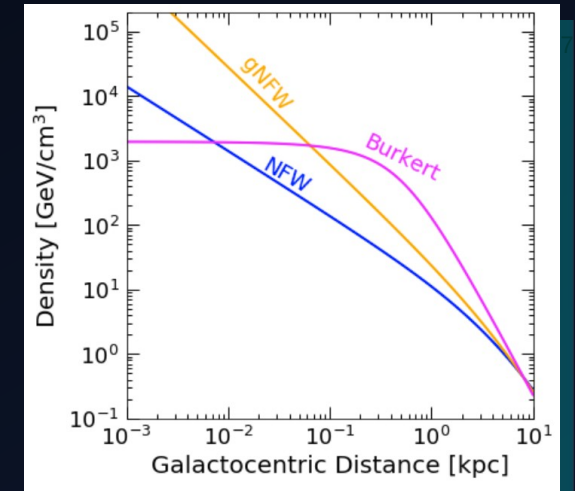
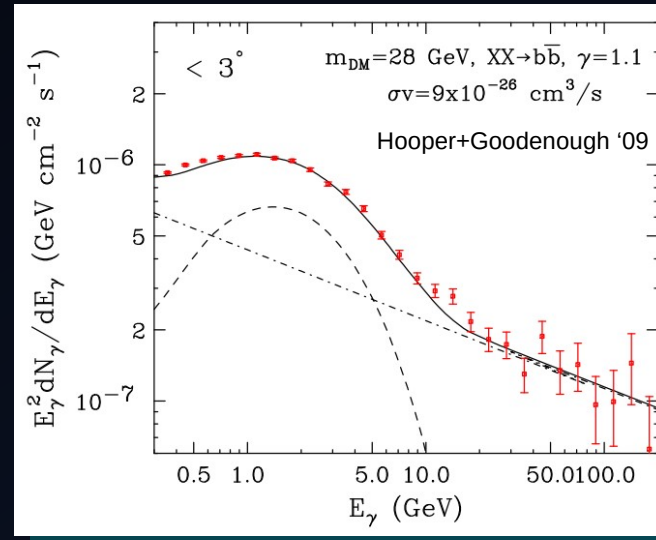
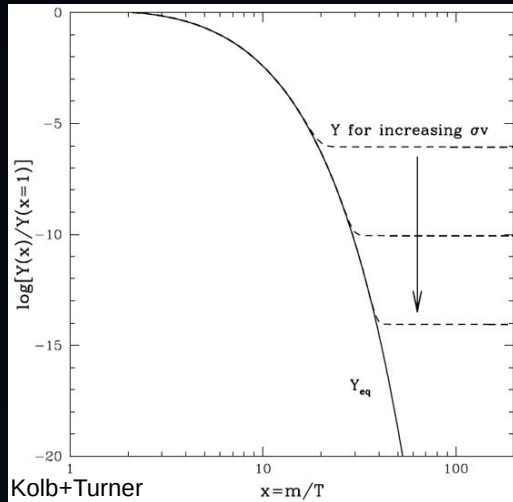
$$\Phi(E, \phi) = \frac{\Gamma}{4\pi m_\chi^a} \frac{dN}{dE} \int \rho[r, (\ell, \phi)]^a d\ell$$



Annihilation cross section

Energy spectrum

"J factor", DM density



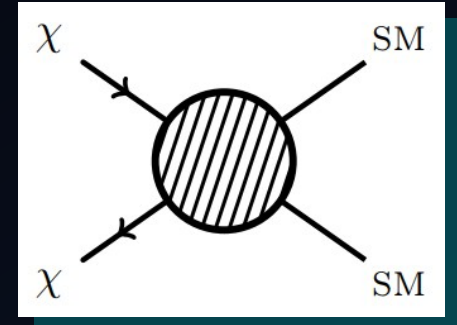
Indirect Detection Ingredients

Particle Physics

Astrophysics

(Neutral particles)

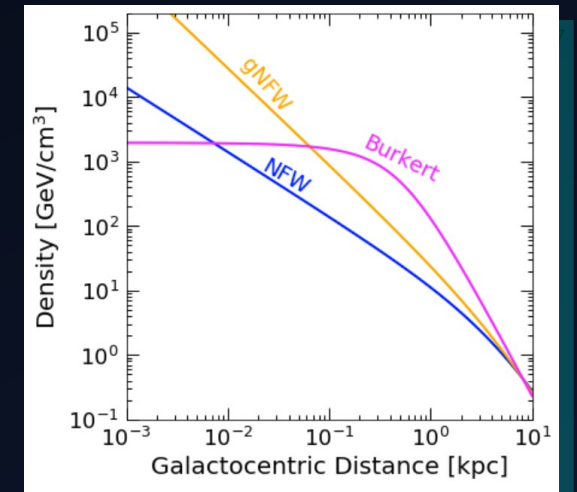
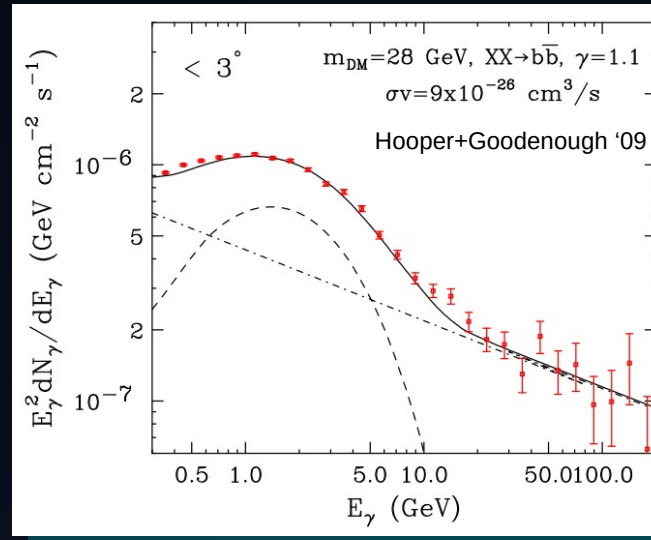
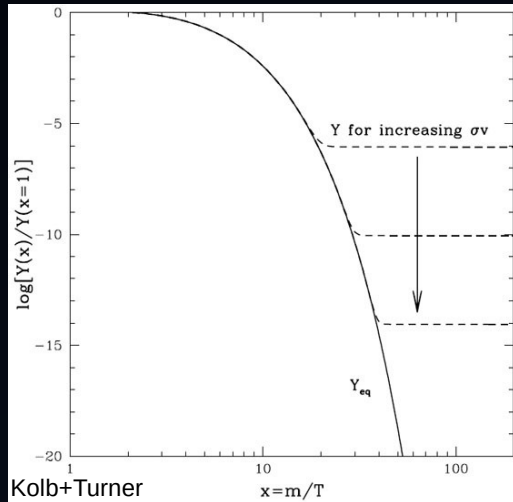
$$\Phi(E, \phi) \frac{\Gamma}{4\pi m_\chi^a} \frac{dN}{dE} \int \rho[r, (\ell, \phi)]^a d\ell$$



Annihilation cross section

Energy spectrum

“J factor”, DM density



Look where this is large!

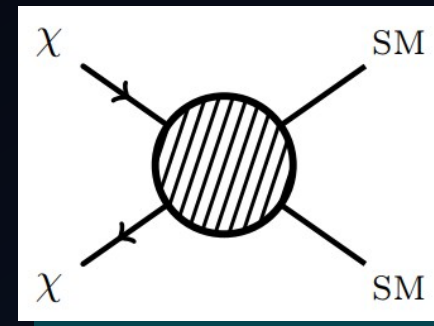
Indirect Detection Ingredients

Particle Physics

Astrophysics

(Neutral particles)

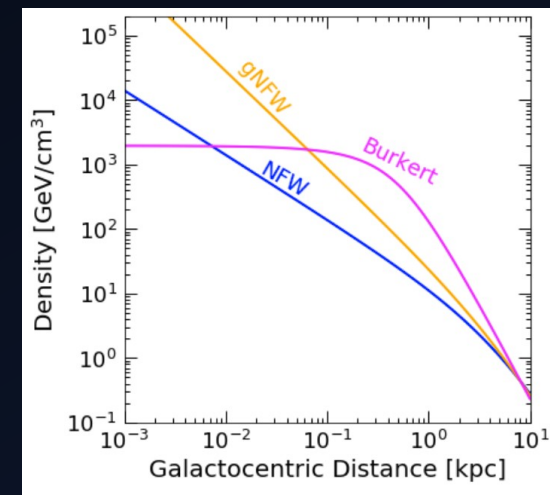
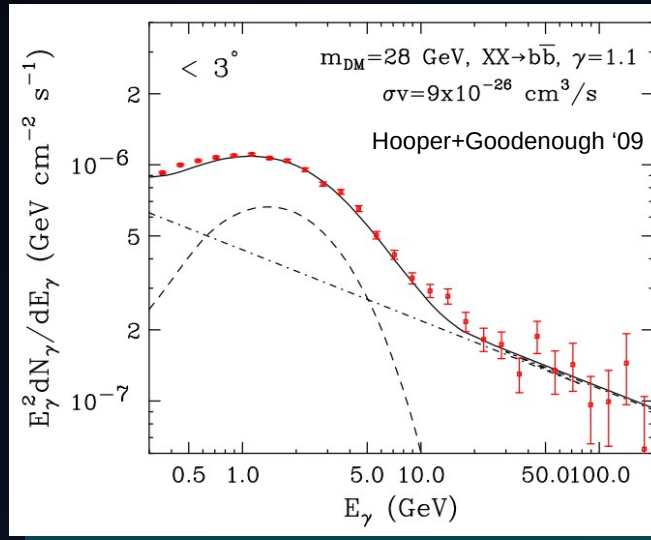
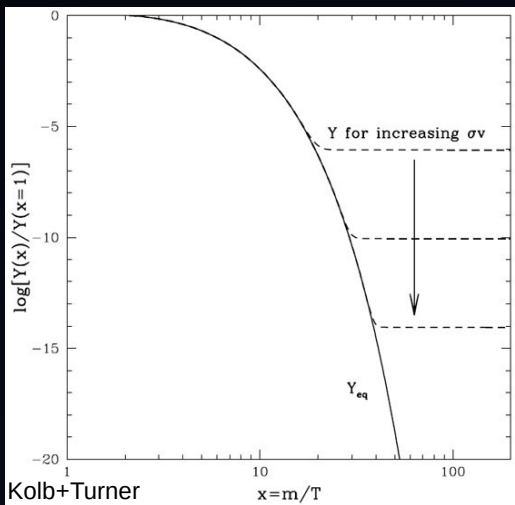
$$\Phi(E, \phi) \frac{\Gamma}{4\pi m_\chi^a} \frac{dN}{dE} \int \rho[r, (\ell, \phi)]^a d\ell$$



Annihilation cross section

Energy spectrum

"J factor", DM density

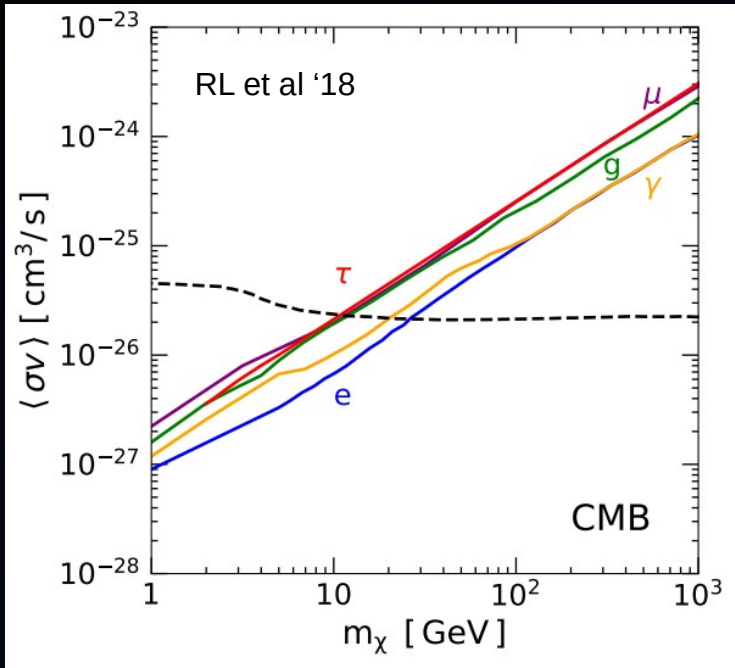


Look where this is large!

...or places with low background!

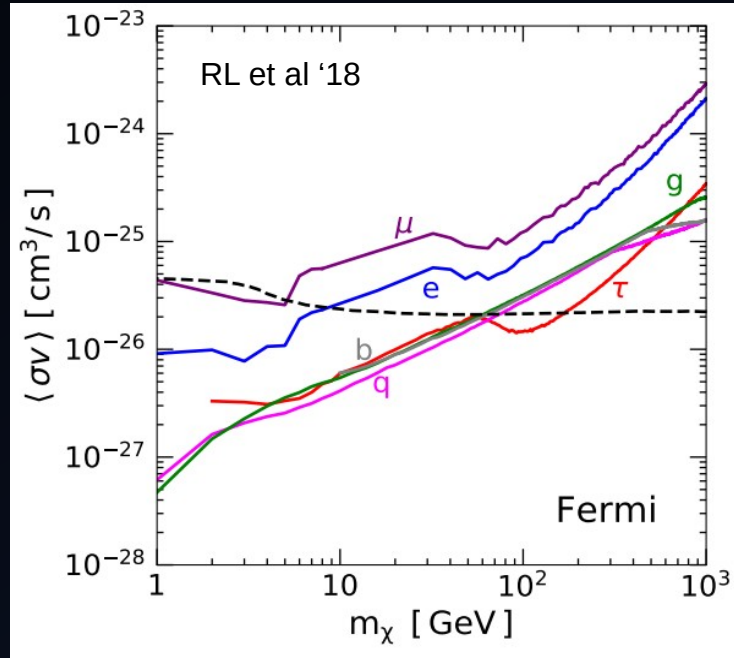
Complementarity: cornering WIMPs

Strongest low mass



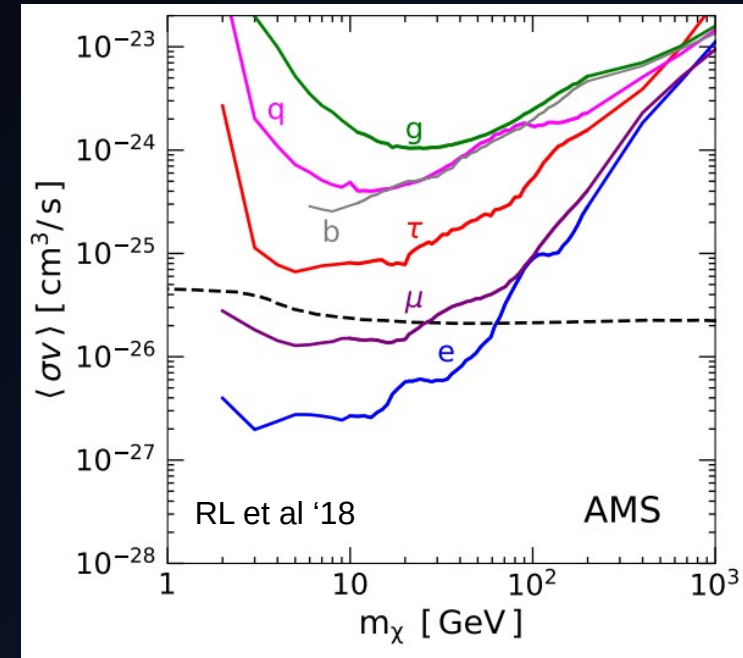
Also see Slatyer '15

Strongest for hadrons



Also see Fermi Collab '16

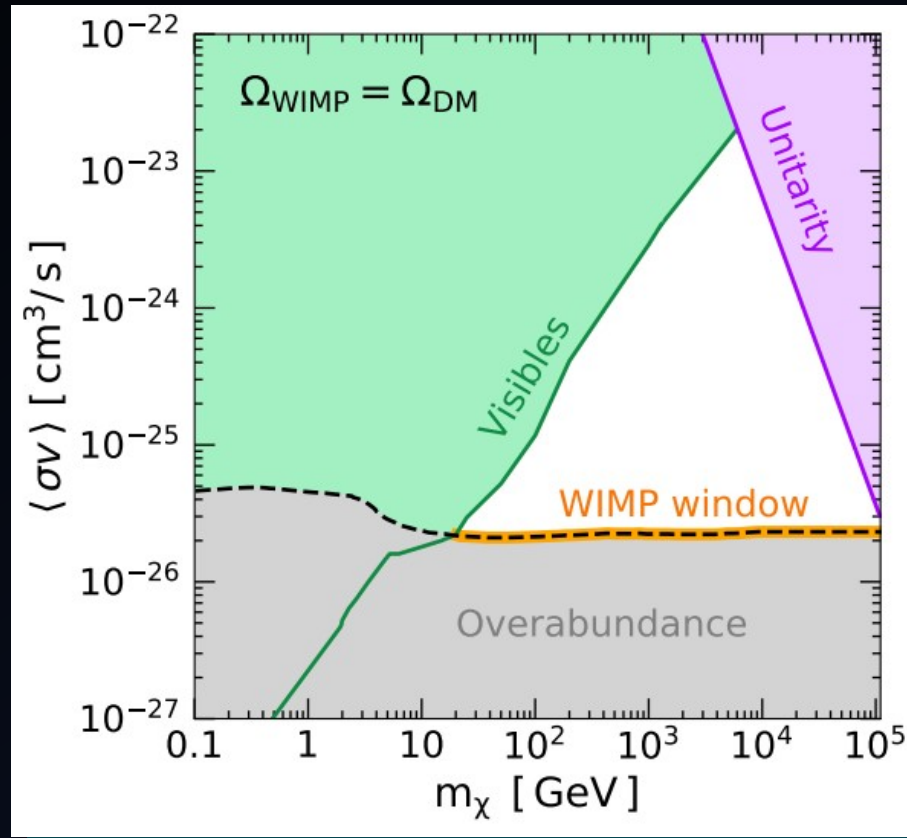
Strongest for leptons



Also see AMS collab '14

(strongest *and* most robust bounds)

Complementarity: cornering WIMPs



RL, Slatyer, Beacom, Ng, '18

WIMP is not dead!

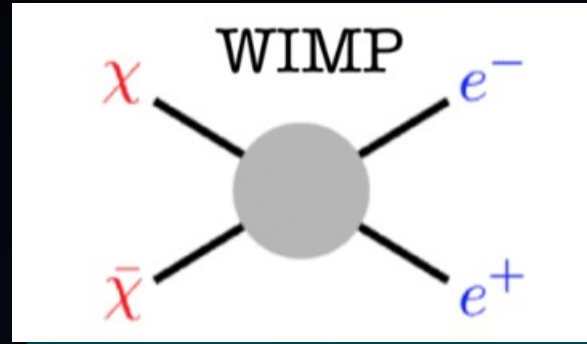
Use all possible final states, combine strongest limits
S-wave $2 \rightarrow 2$ thermal DM to visible states: mass greater than ~ 20 GeV

Still need to push through this window

OTHER TYPES OF FREEZEOUT

Range of freezeout options!

Zeldovich, Lee,
Weinberg, Steigman,
Turner,...+

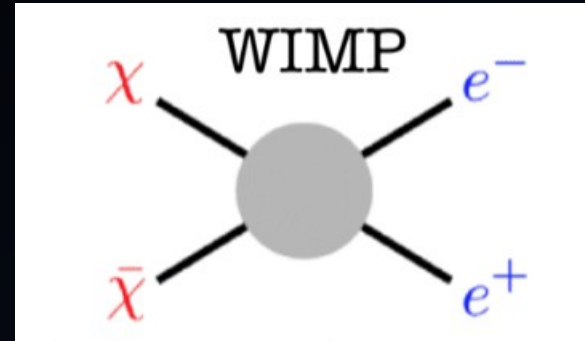


$$\Gamma_{\text{DM}} = \langle \sigma v_{\text{rel.}} \rangle n_{\text{DM}} > H(T)$$

$$\Omega_{\text{DM}} h^2 \approx \frac{0.12}{\langle \sigma v_{\text{rel.}} \rangle [25 \text{TeV}]^2}$$

Range of freezeout options!

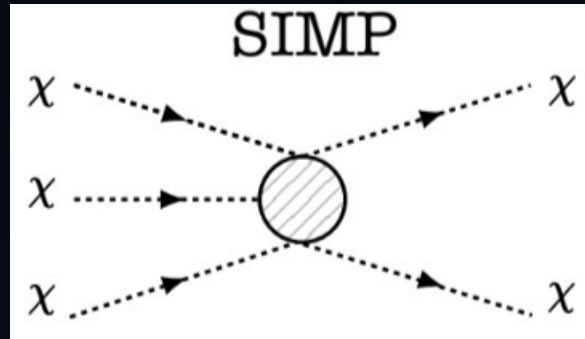
Zeldovich, Lee,
Weinberg, Steigman,
Turner,...+



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Hochberg, Kuflik,
Volanksy, Wacker, 2014

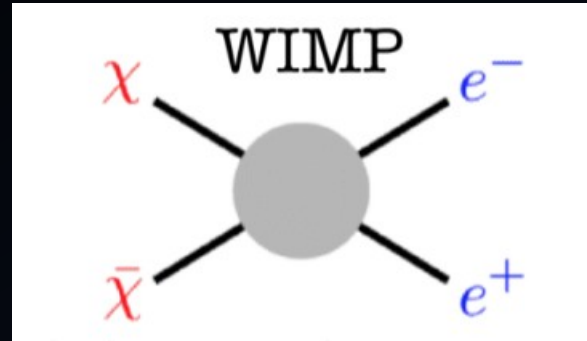


$$\Gamma_{\text{DM}} = \langle \sigma_{3 \rightarrow 2} v_{\text{rel.}}^2 \rangle n_{\text{DM}}^2 > H(T)$$

$$\Omega_{\text{DM}} h^2 \approx \left(\frac{\text{MeV}}{m_{\text{DM}}} \right) \frac{0.12}{\sqrt{\langle \sigma_{3 \rightarrow 2} v_{\text{rel.}}^2 \rangle [3 \text{MeV}]^5}}$$

Range of freezeout options!

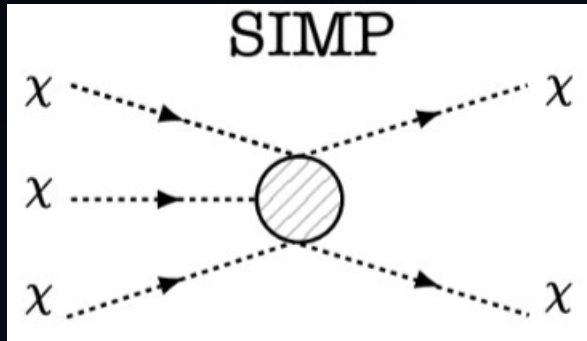
Zeldovich, Lee,
Weinberg, Steigman,
Turner,...+



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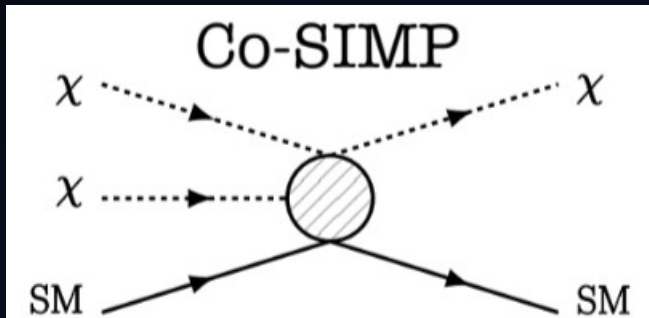
Hochberg, Kuflik,
Volanksy, Wacker, 2014



$$\Gamma_{\text{DM}} = \langle \sigma_{3 \rightarrow 2} v_{\text{rel.}}^2 \rangle n_{\text{DM}}^2 > H(T)$$

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Beacom, Smirnov, 2020



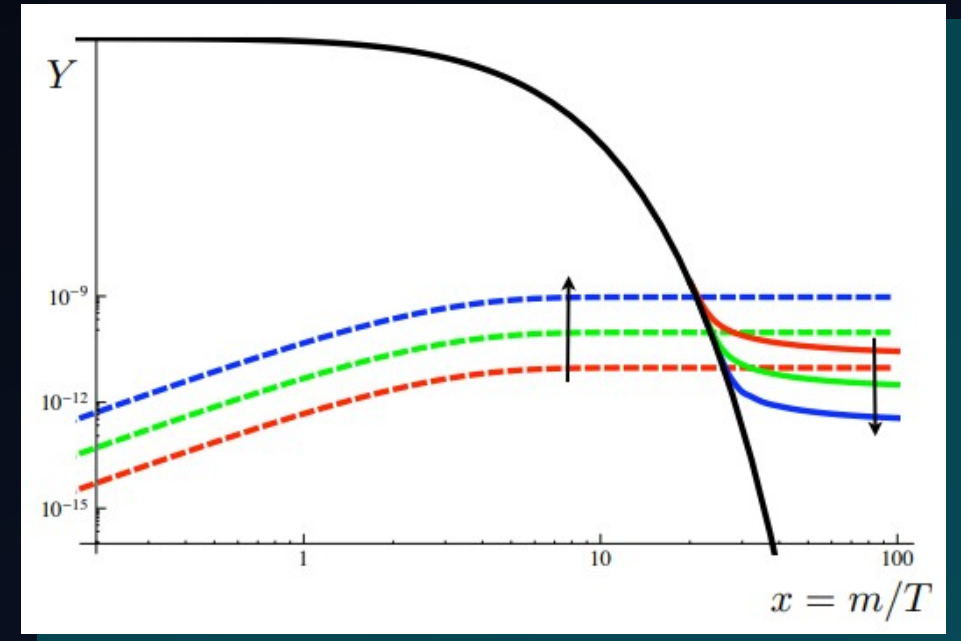
$$\Gamma_{\text{DM}} = \langle \sigma_{3 \rightarrow 2} v_{\text{rel.}}^2 \rangle n_{\text{DM}} n_{\text{SM}} > H(T)$$

$$\Omega_{\text{DM}} h^2 \approx \left(\frac{\text{MeV}}{m_{\text{DM}}} \right)^3 \frac{0.12}{\langle \sigma_{3 \rightarrow 2} v_{\text{rel.}}^2 \rangle [100 \text{ MeV}]^5}$$

FREEZEOUT? FREEZE IN!

Dark matter freeze-in

- For freeze-out, the universe has a large population of DM that was in thermal equilibrium with the bath, that is depleted as the Universe cools
- For freeze-in, the universe starts with little DM and it never reaches thermal equilibrium with the bath
- At temperatures lower than the DM mass, the bath no longer has enough energy to produce it. Then, the DM is “frozen-in”

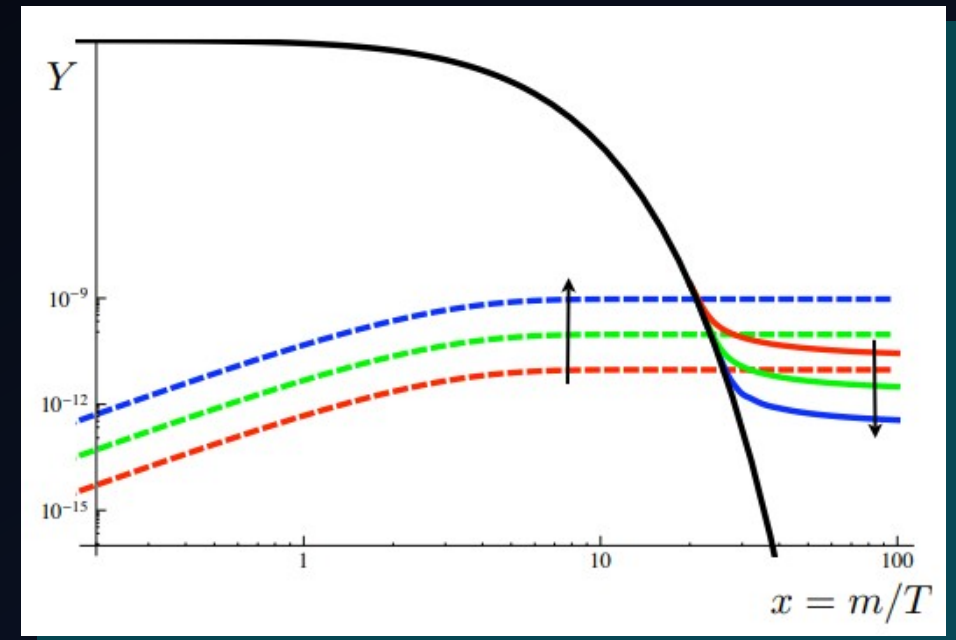


Arrows point in direction of increased DM-bath coupling

Hall, Jedamzik,
March-Russell, West
2009

Dark matter freeze-in

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- At temperatures lower than the DM mass, the bath no longer has enough energy to produce it. Then, the DM is “frozen-in”
- Still gets us right abundance, but applies to models with much smaller couplings -- “Feebly Interacting Massive Particles”, or FIMPs

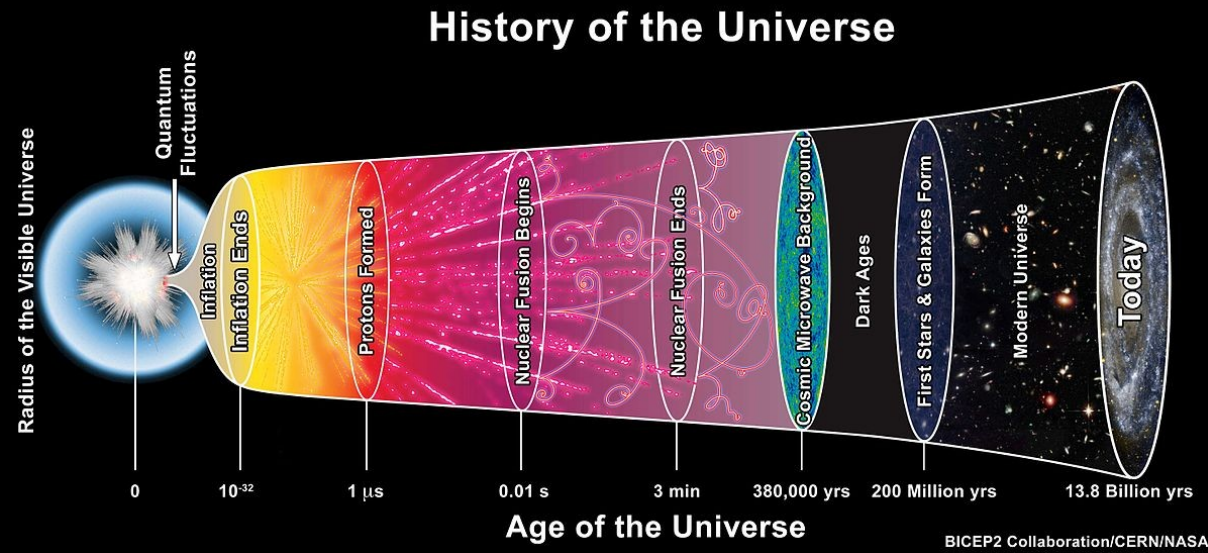


Arrows point in direction of increased DM-bath coupling

Hall, Jedamzik,
March-Russell, West
2009

Uncertainties for the early Universe?

- WIMP freezes out before BBN and CMB
 - A lot could happen between these times!
- Some particle could decay into DM
- Some particle could decay into visible matter, diluting the original amount of DM



- Maybe physics was just different to what we expect! Hard to know.
 - Finding DM may give us a new window further back into the Universe

AXIONS

(see also Jure Zupan's earlier lecture)

Strong CP problem

- If interactions are allowed by all symmetries, they generally are present
- One term is suspiciously silent

$$\mathcal{L}_\theta = \frac{\theta}{16\pi^2} G_{\mu\nu} \tilde{G}_{\mu\nu}$$

- Strange also because symmetry between matter and antimatter is broken in the weak interactions (CP violated), but not in strong interactions
- Term should induce neutron electric dipole moment, but experimentally it is measured to be at least **~10 ORDERS OF MAGNITUDE below naive expectations ????**

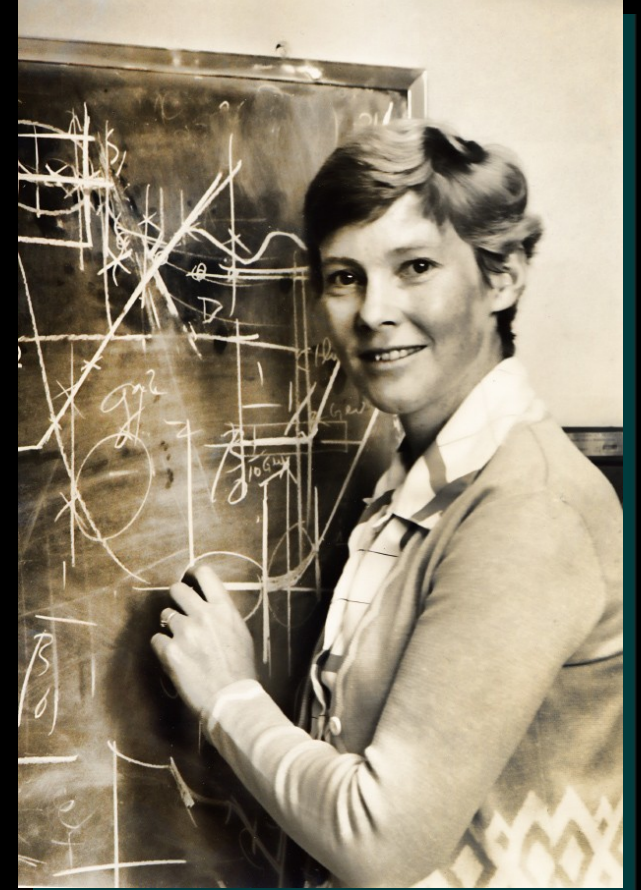
ENTER PECCEI + QUINN



Image: Quanta Magazine

Peccei-Quinn Symmetry

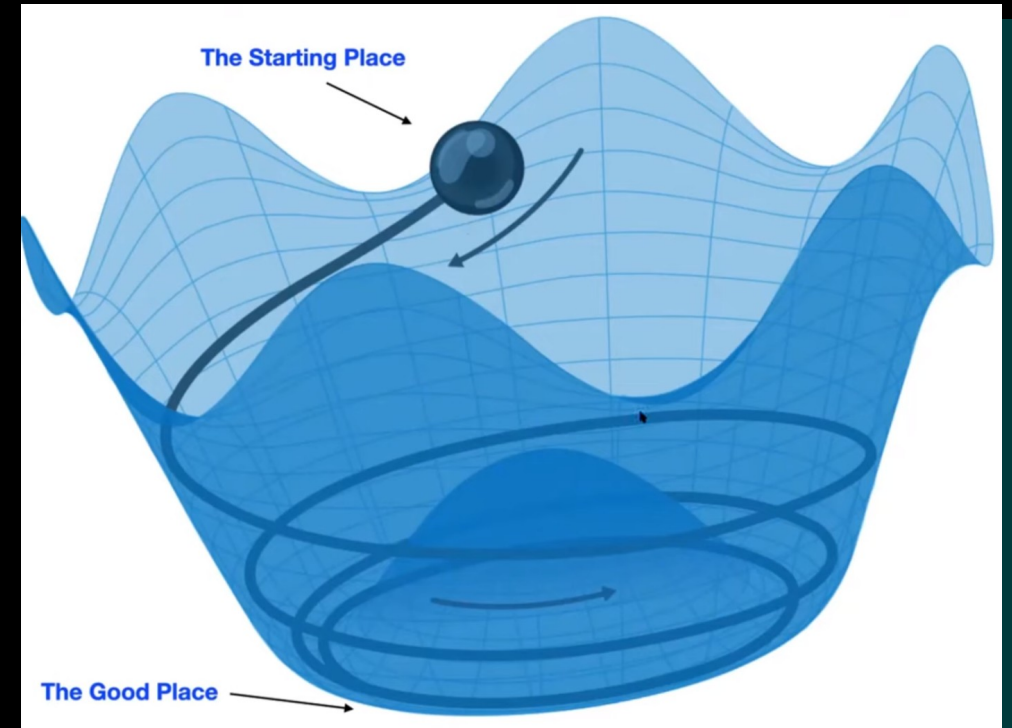
- If interactions are allowed by all symmetries, they generally are present
- So: Invoke new symmetry, Peccei-Quinn (PQ) symmetry, which is a minimal extension of the SM, which is spontaneously broken



Helen Quinn in her SLAC office, ~1977
Source: Helen Quinn

Dynamical field

- Promote the unwanted coupling (theta) to a dynamical field that relaxes to zero: not fixed for all time, but instead changes over time
- Value of the field is the co-efficient of the "unwanted" interaction
- Energetically preferred coupling is small/zero (so over time of the Universe it evolves to zero)
- Quanta of this field is the QCD axion



Peccei, Quinn 1977
Wilczek 1978, Weinberg 1978
Vafa, Witten 1984

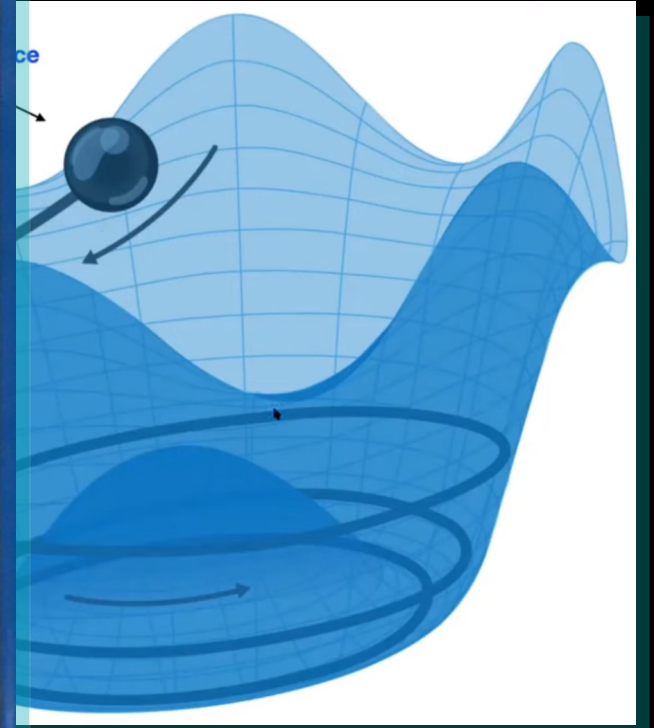
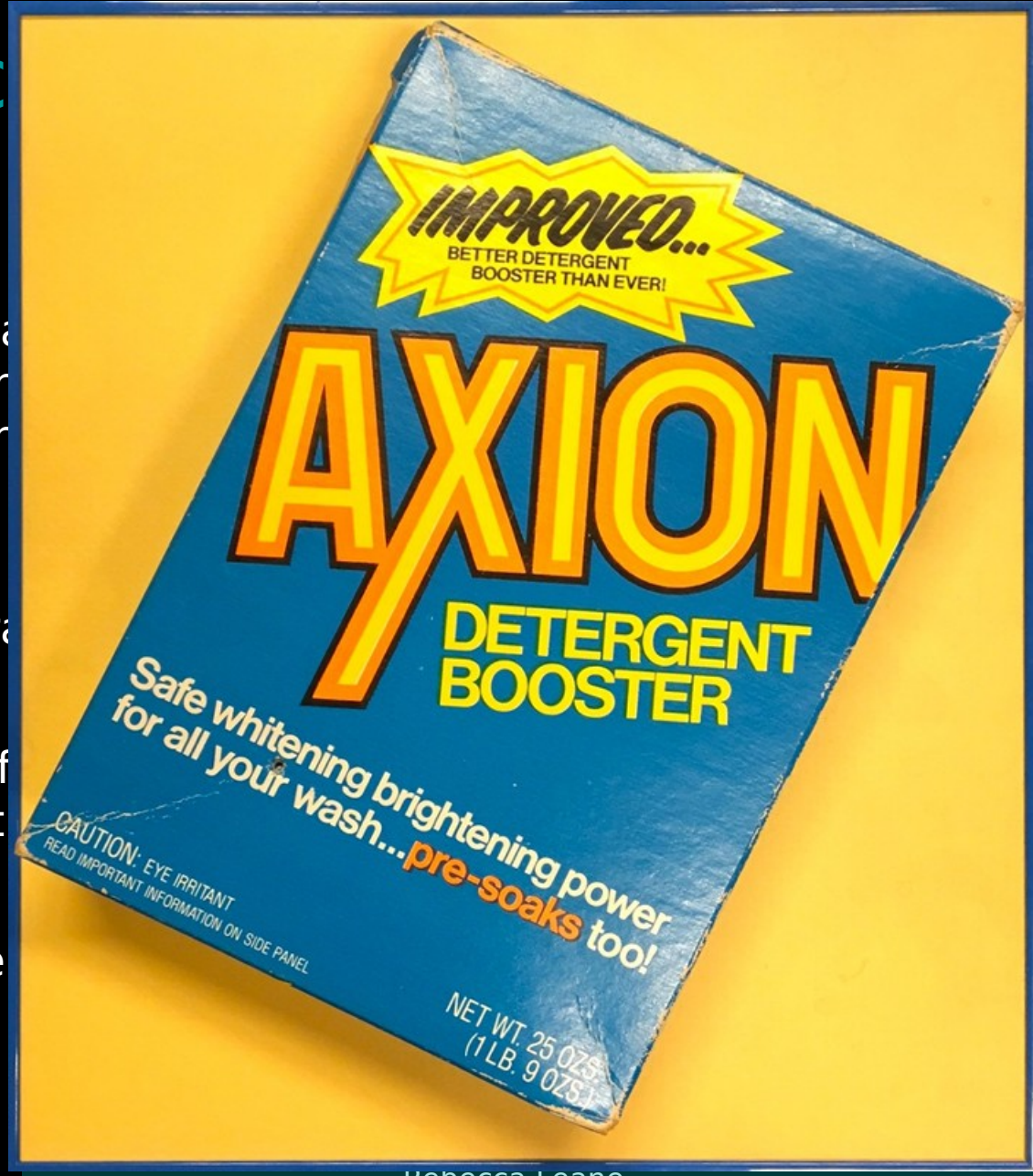
$$m_a = 5.7 \mu\text{eV} \left(\frac{10^{12} \text{GeV}}{f_a} \right)$$

Rebecca Leane

Wilczek

Dynamic

- Promote the unwanted dynamical field that exists for all time, but in a way that is not wanted
- Value of the field is "unwanted" interaction
- Energetically preferred (so over time of time)
- Quanta of this field

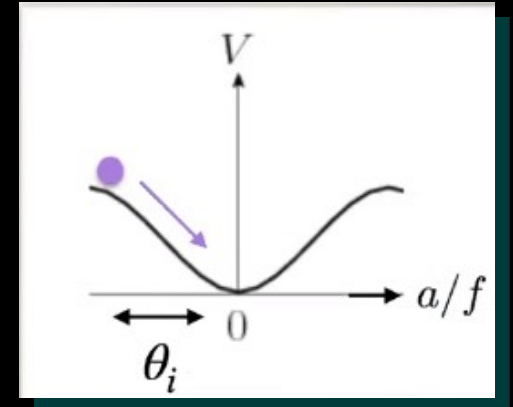


Rebecca Leane

Wilczek

Can be dark matter

- Production occurs in the early universe, misalignment mechanism
 - Relic abundance ok for small masses, $\sim 10^{-8} - 10^{-3}$ eV
- Generally want mass less than about 20 eV for stability
- Can act as cold dark matter!



Zupan's slide



Increasingly popular, and can be DM!

A New Light Boson?

Steven Weinberg (Harvard U.)

Dec, 1977

12 pages

Published in: *Phys.Rev.Lett.* 40 (1978) 223-226

DOI: [10.1103/PhysRevLett.40.223](https://doi.org/10.1103/PhysRevLett.40.223)

Report number: HUTP-77/A074

View in: [OSTI Information Bridge Server](#), [ADS Abstract Service](#), [KEK scanned document](#)

[pdf](#) [cite](#) [claim](#)

↻ 4,670 citations

Citations per year



Frank Wilczek (Columbia U. and Princeton, Inst. Advanced Study)

Nov, 1977

12 pages

Published in: *Phys.Rev.Lett.* 40 (1978) 279-282

DOI: [10.1103/PhysRevLett.40.279](https://doi.org/10.1103/PhysRevLett.40.279)

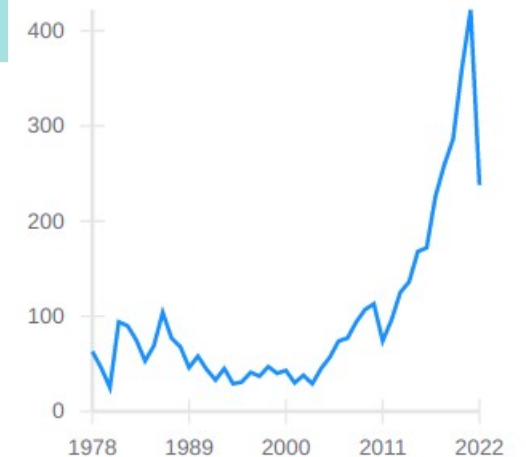
Report number: Print-77-0939 (COLUMBIA)

View in: [OSTI Information Bridge Server](#), [ADS Abstract Service](#), [KEK scanned document](#)

[pdf](#) [cite](#) [claim](#)

↻ 4,484 citations

Citations per year



Testable interactions

- Can couple to gluons, EM, leptons, quarks

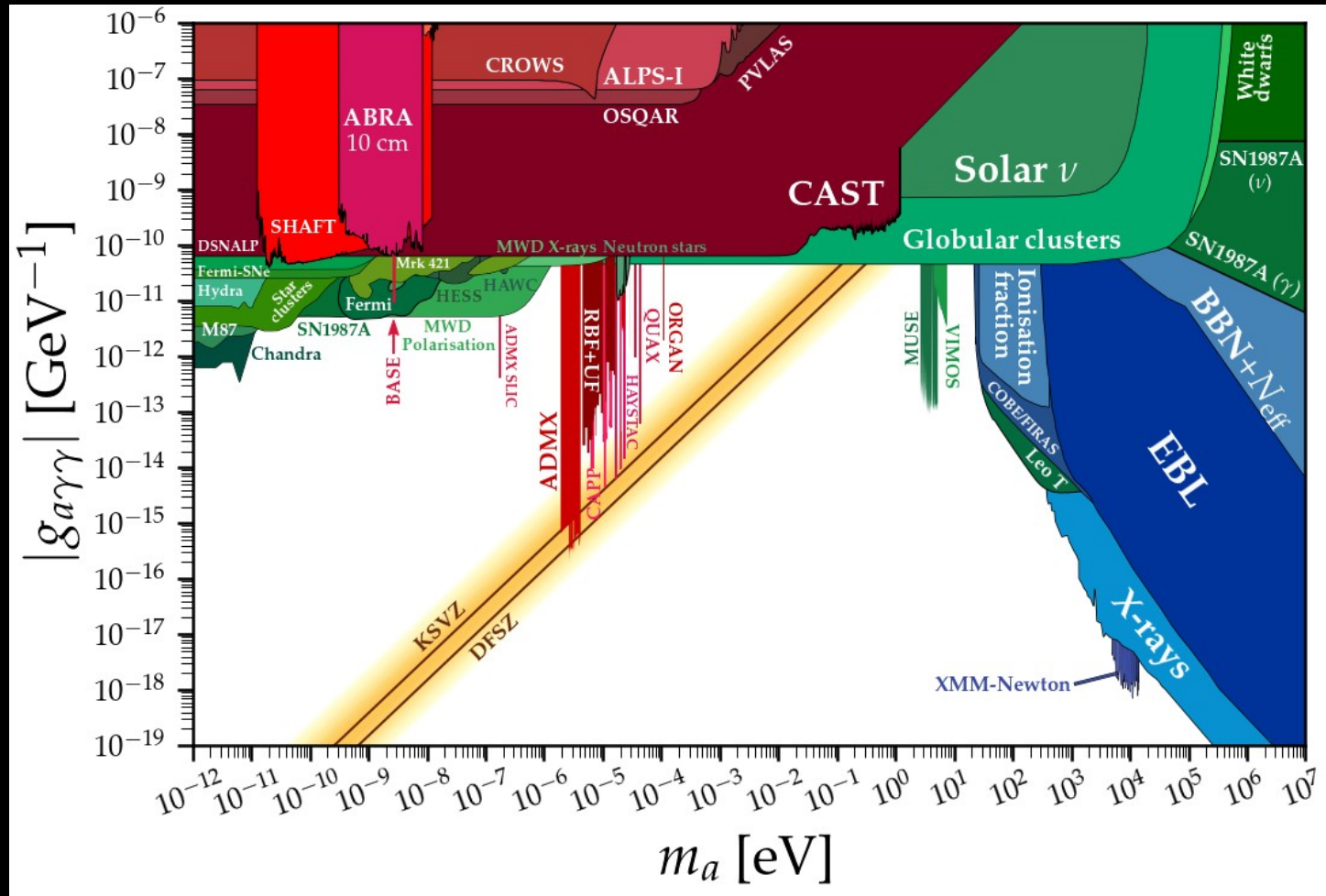
$$\mathcal{L}_{\text{int}} \sim -\frac{a}{F} (c_G \alpha_s G_{\mu\nu} \tilde{G}^{\mu\nu} + c_\gamma \alpha F_{\mu\nu} \tilde{F}^{\mu\nu} + d_q \sum_q m_q \bar{q} \gamma_5 q + d_l \sum_l m_l \bar{l} \gamma_5 l + \dots)$$

- Modifies Maxwell's equations:

$$\begin{aligned}\nabla \cdot E &= -\kappa \nabla a \cdot B \\ \nabla \times E &= -\frac{\partial B}{\partial t} \\ \nabla \cdot B &= 0 \\ \nabla \times B &= \frac{\partial E}{\partial t} + \kappa (\dot{a} B + \nabla a \times E)\end{aligned}$$

Some interactions can also apply to axion-like particles (ALPs)

Axion limits

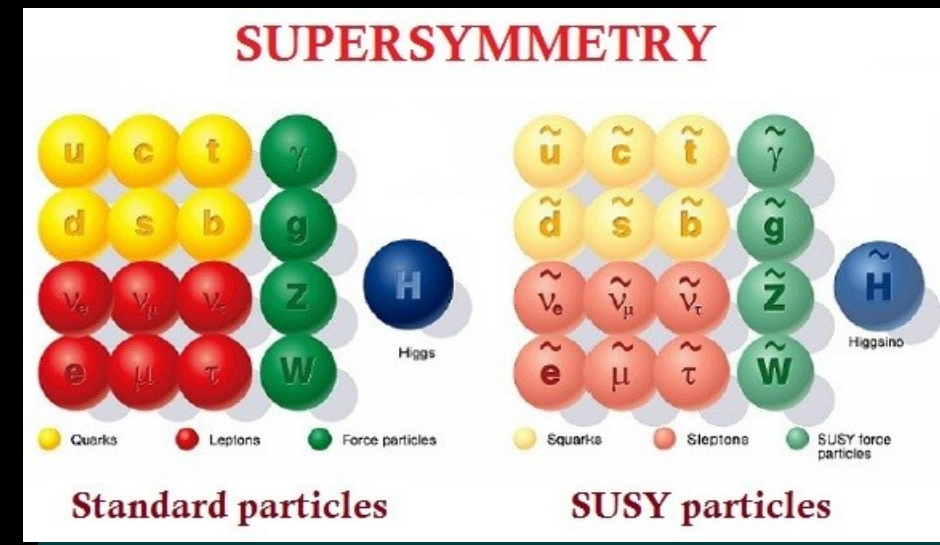


<https://cajohare.github.io/AxionLimits/>

COMPLETE DM MODELS

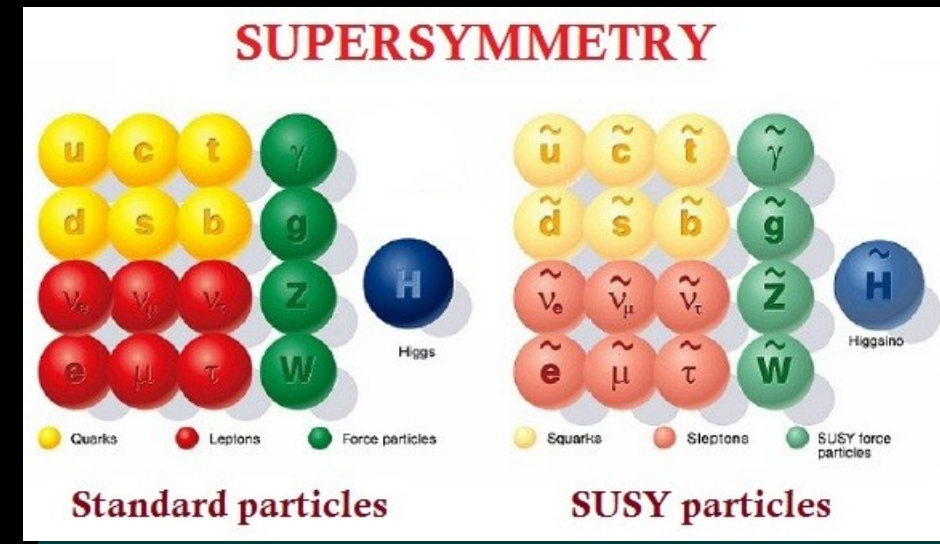
Complete DM models

- **Axions:**
naturally appear in string theory, GUTs
- **Supersymmetry:**
Prevailing new physics model for a long time, naturally gave us:
 - Dark matter candidate (avoid proton decay enforcing R-parity, so superpartners only couple in pairs to the SM \rightarrow LSP)
 - Hierarchy problem solution, unification...



Complete DM models

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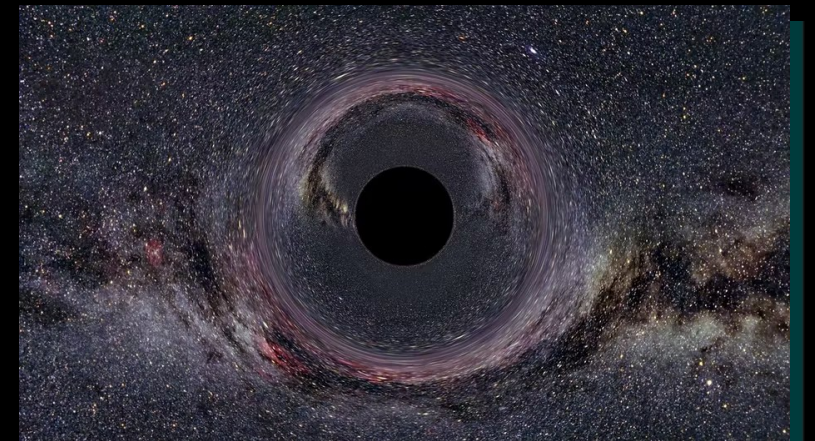
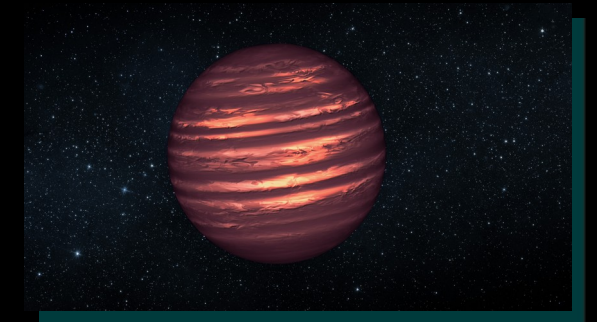


Complete models can guide us where us to look for other new physics
(see unification talks later this week)

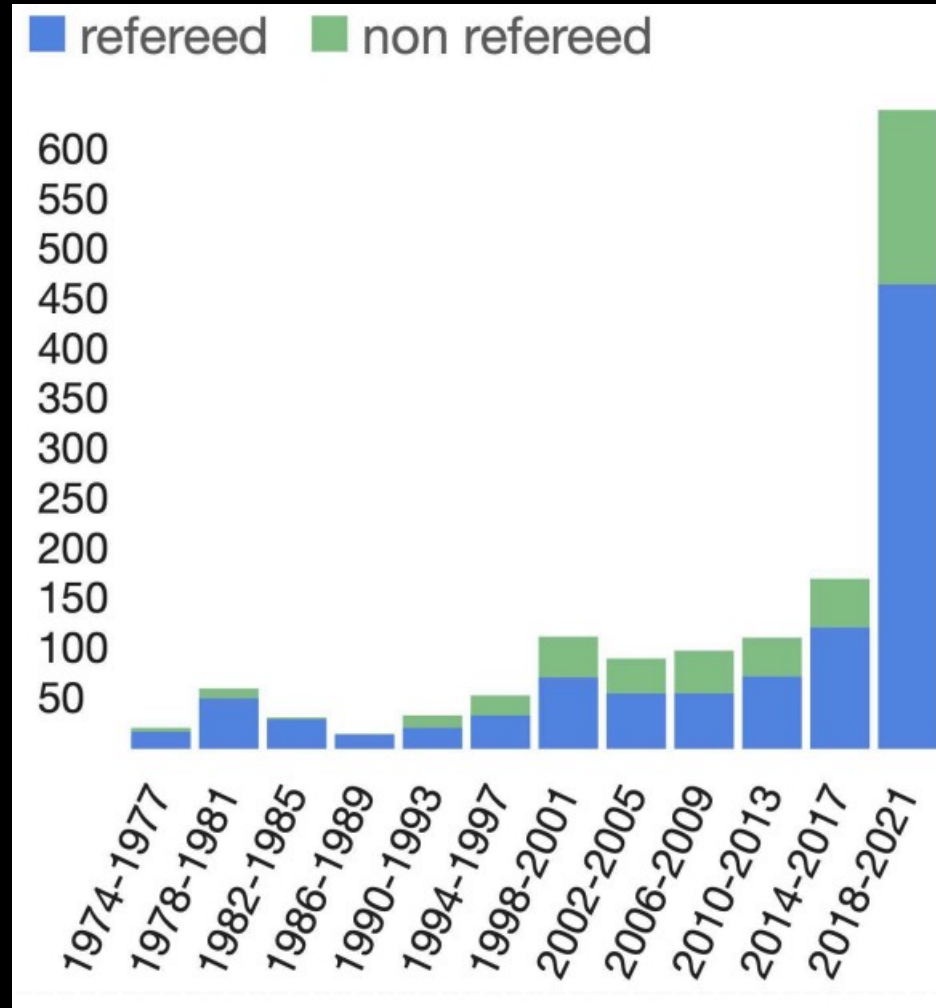
MACHOS: COMPOSITE DARK MATTER

MACHOs

- Compact objects like black holes, neutron stars, brown dwarfs...
- Originally, was just thought could be baryonic objects
 - Hard to reconcile with observations:
 - CMB / BBN
 - Microlensing surveys
- Increasingly popular candidate: primordial black holes



Primordial Black Hole Popularity



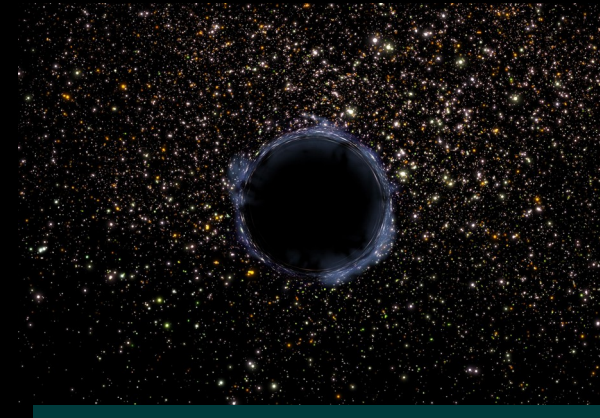
Carr + Kuhnel, 2022

Primordial Black Holes

- First considered in 1966 by Zel'dovich and Novikov, in more detail by Hawking and Carr in the early 1970s
- What if DM is just black holes that formed before the epoch of BBN, AND with masses below the sensitivity range of microlensing surveys?
- Viable option, though we need (1) some production mechanism, and (2) some explanation of the current constraints from again CMB and lensing, among other things



Primordial Black Holes: Production



- Cosmological density time t after the Big Bang is

$$\rho \sim 1/(G t^2)$$

- Density needed for volume with mass M to fall within Schwarzschild radius is

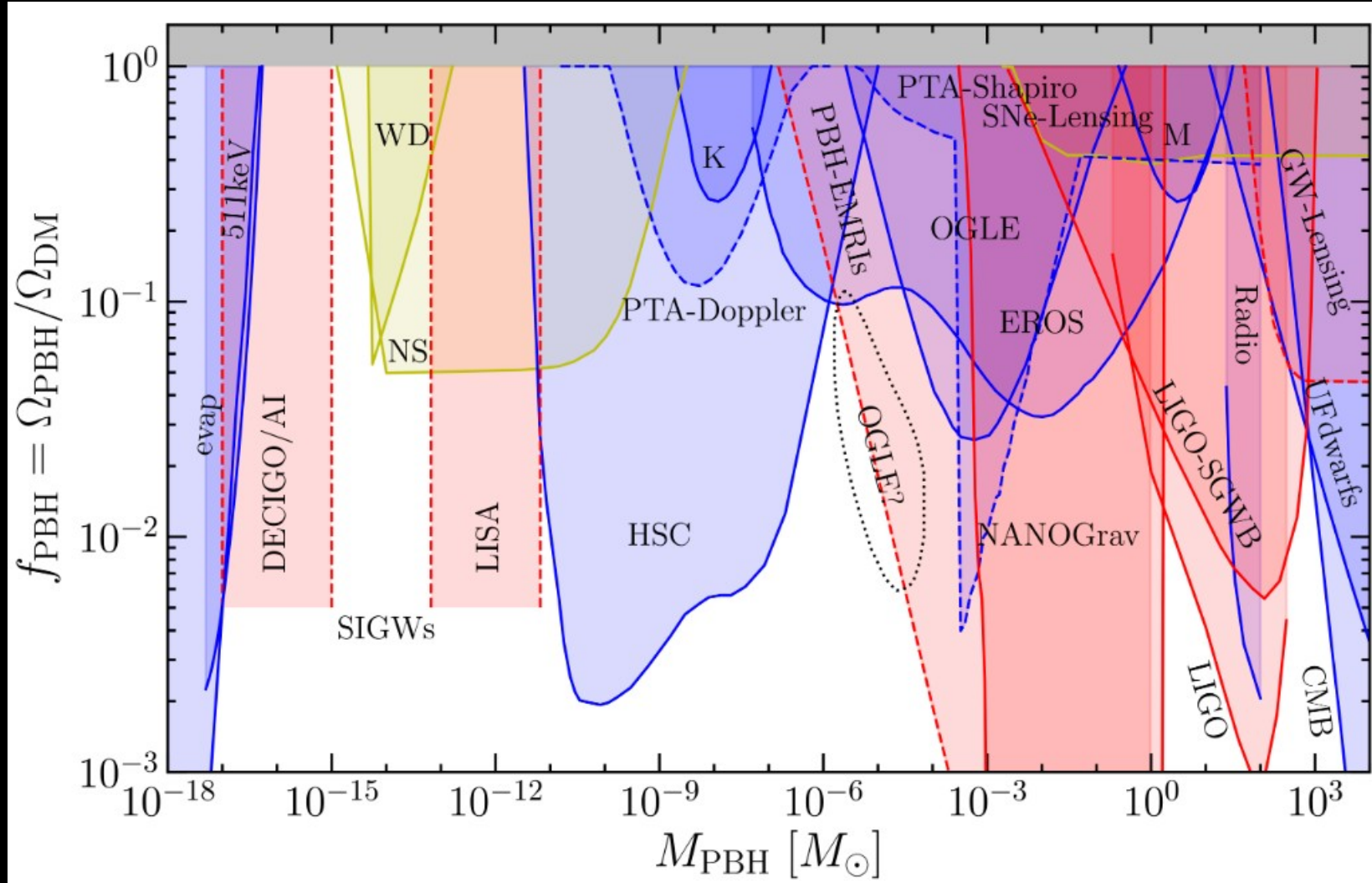
$$\rho \sim c^6/(G^3 M^2)$$

- Tells us that around the cosmological horizon, mass of PBHs would be

$$M \sim \frac{c^3 t}{G} \sim 10^{15} \left(\frac{t}{10^{-23} \text{ s}} \right) \text{ g}$$

- Therefore, depending on formation time, could span huge mass range!

Primordial Black Holes: Constraints



Outlook for dark matter

- **Dark matter exists!** Diverse range of evidence across many length scales
 - Finding its nature is a key goal of our community
- **Hard to know the right dark matter theory direction.** No clear excesses to explain, no clear “correct” model.
- **Experiment and theory have historically informed each other,** and will continue to do so. Extensive new technological advances and new instruments on the horizon.
- **Dark matter could be around the corner!**

